

Interactions

Take several fields ϕ_i, ψ_i, F_i

$$\delta\phi_i = \epsilon\phi_i \quad \delta\phi_i^* = \epsilon^*\psi_i^\dagger$$

$$\delta\psi = -i(\sigma^\mu \epsilon^\dagger)\partial_\mu\psi + \epsilon F_i$$

$$\delta\psi^\dagger = i(\epsilon\sigma^\mu)\partial_\mu\psi_i^* + \epsilon^\dagger F_i^*$$

$$\delta F_i = -i\epsilon^\dagger\bar{\sigma}^\mu\partial_\mu\psi_i$$

$$\delta F_i^* = i\partial_\mu\psi_i^\dagger\sigma^\mu\epsilon$$

$$\mathcal{L} = -\partial_\mu\phi_i^*\partial^\mu\phi_i + i\bar{\psi}_i\bar{\sigma}^\mu\partial_\mu\psi_i + F_i^*F_i$$

Can I add interactions? dimensions are

ϕ	: 1
ψ	: 3/2
F	: 2

• can have $\alpha_{ij}^{\mu\nu}F_iF_j, F_i^*F_j^*$ ($F_iF_j^*$ would interfere with $\psi_i^\dagger\psi_j$)

$$W_{ij}^{\mu\nu}F_i$$

$$W_{ij}^{\mu\nu}\psi_i\psi_j$$

+ hermitian conjugates

remember that $\psi_i^\dagger\psi_j$ cannot be made Lorentz invariant; only via $\psi_i^\dagger\bar{\sigma}^\mu\psi_j$
and adding ∂^μ is the term we already have in the kinetic terms

Interactions of a chiral multiplet

$$\mathcal{L}_{\text{free}} = -\partial_\mu \phi_i \partial^\mu \phi_i + i \psi_i^\dagger \bar{\sigma}^\mu \partial_\mu \psi_i + F_i^\dagger F_i$$

$$\delta \phi_i = \epsilon \psi_i$$

$$\delta \phi_i^\dagger = \epsilon^\dagger \psi_i^\dagger$$

$$\delta F_i = -i \epsilon^\dagger \bar{\sigma}^\mu \partial_\mu \psi_i$$

$$\delta F_i^\dagger = i \partial_\mu \psi_i^\dagger \bar{\sigma}^\mu \epsilon$$

$$\delta \psi_i = (-i \sigma^\mu \epsilon^\dagger) \partial_\mu \phi_i + \epsilon F_i$$

$$U(\phi_i, \phi_j^\dagger) \rightarrow \underbrace{\frac{\partial U}{\partial \phi_i}}_{\phi \phi^\dagger \epsilon \psi_i} + \dots$$

cannot be produced by the others

$$\left. \begin{array}{l} F_i^\dagger F_j \\ F_i^\dagger F_j^\dagger \end{array} \right\} \rightarrow \left. \begin{array}{l} F_i \\ F_i \end{array} \right\} (-i) \epsilon^\dagger \bar{\sigma}^\mu \partial_\mu \psi_i$$

cannot be produced by the other

$$\begin{aligned} & F_i^\dagger W^i(\phi, \phi^\dagger) \rightarrow -i \epsilon^\dagger \bar{\sigma}^\mu \partial_\mu \psi_i W^i(\phi, \phi^\dagger) + F_i^\dagger \epsilon \psi_j \frac{\partial W^i}{\partial \phi_j} \\ & + \text{h.c.} \\ & \psi_i \psi_j W^{ij}(\phi, \phi^\dagger) \rightarrow -i 2 \psi_i \sigma^\mu \epsilon^\dagger \partial_\mu \phi_j W^{ij}(\phi, \phi^\dagger) \\ & + \psi_i \psi_j \frac{\partial W^{ij}}{\partial \phi_k} \epsilon \psi_k + \psi_i \psi_j \frac{\partial W^{ij}}{\partial \phi^* k} \epsilon^\dagger \psi_k^\dagger \\ & + \text{h.c.} \\ & + 2 \psi_i \epsilon F_j W^{ij}(\phi, \phi^\dagger) \end{aligned}$$

① + fermions $\frac{\partial W^{ij}}{\partial \phi^k}$ must be fully symmetric in ijk

② $\frac{\partial W^{ij}}{\partial \phi^* k} = 0$

$$U(\phi_i, \phi_j^*) \Rightarrow \frac{\partial U}{\partial \phi_i} \varepsilon \psi_i + \frac{\partial U}{\partial \phi_j^*} \varepsilon^+ \psi_i^+$$

$F_i F_j^*$: mixes with the kinetic terms

$$F_i F_j x^{ij} \Rightarrow 2 x^{ij} F_j (-i) \varepsilon^+ \bar{\sigma}^\mu \partial_\mu \psi$$

$$F_i W^i(\phi_j, \phi_k^*) \Rightarrow W^i (-i) \varepsilon^+ \bar{\sigma}^\mu \partial_\mu \psi_i + F_i \frac{\partial W^i}{\partial \phi_j} \varepsilon \psi_j + F_i \frac{\partial W^i}{\partial \phi_j^*} \varepsilon^+ \psi_j^+$$

$$\psi_i \psi_j W^{ij}(\phi, \phi^*) \Rightarrow 2$$

Gauge supermultiplet

	A_μ	λ	D
on-shell	2	2	0
off-shell	3	4	1

dim	A^a	
	λ^a	$3/2$
	D^a	2

$$\delta A_\mu^a = \left(\epsilon^\dagger \bar{\sigma}_\mu^a + \lambda^a \bar{\sigma}_\mu^{a\dagger} \epsilon \right) \times \frac{1}{\sqrt{2}}$$

$$\delta \lambda^a = \frac{i}{2\sqrt{2}} (\sigma^\mu \bar{\sigma}^\nu \epsilon) F_{\mu\nu}^a + \frac{1}{\sqrt{2}} \epsilon D^a$$

$$\delta D^a = \frac{i}{\sqrt{2}} \left(-\epsilon^\dagger \bar{\sigma}^\mu \not{\partial} \lambda^a + \not{\partial}_\mu \lambda^{a\dagger} \bar{\sigma}^\mu \epsilon \right)$$

overall factor fixed such that $[\delta_{\epsilon_1}, \delta_{\epsilon_2}] X^a = \left(-\frac{\epsilon_1^\dagger \sigma^\mu \epsilon_2^\dagger + \epsilon_2^\dagger \sigma^\mu \epsilon_1^\dagger}{2} \right) i \not{\partial} X^a$

Note: covariant derivatives

• \Rightarrow nonlinear

• can add (even more) auxiliary fields to make it linear

Can make $\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} + i \lambda^a \bar{\sigma}^\mu \not{\partial}_\mu \lambda^a + \frac{1}{2} D^a D^a$

supersymmetric in this way.

Couple gauge & chiral multiplets

$$\Rightarrow \mathcal{L}_i = -i \epsilon^\mu \bar{\sigma}^\nu \nabla_\nu \psi_i + \sqrt{2} g (T^a \phi)_i \lambda^{a\dagger}$$

& we need extra interactions with $\lambda^a D^a$ as well

$$-\sqrt{2} g (\phi^\dagger T^a \psi) \lambda^a - \sqrt{2} g \lambda^{a\dagger} (\psi^\dagger T^a \phi) + g (\phi^\dagger T^a \phi) D^a$$

are dim ≤ 4

$W^i(T^a \phi)_i$ must be zero \Rightarrow also needed for gauge invariance

\Rightarrow scalar interactions fixed!