

PARTIALLY QUENCHED CHIRAL PERTURBATION THEORY AT TWO LOOPS

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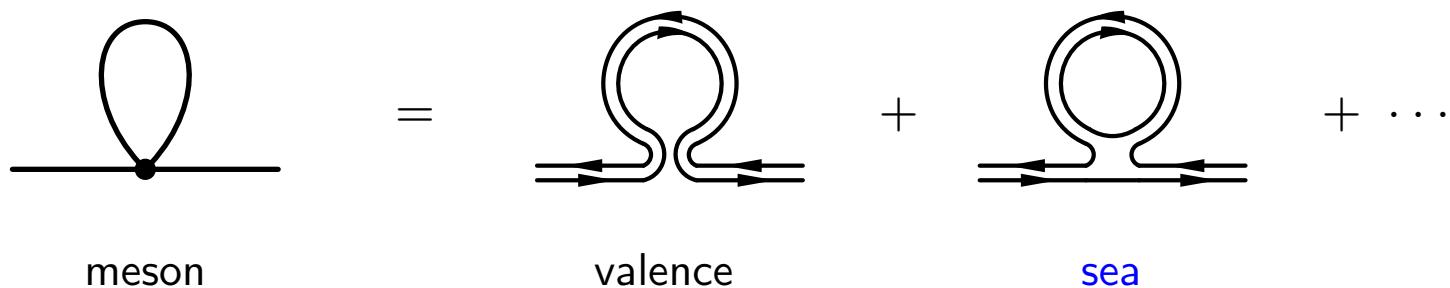
– Timo Lähde – Partikeldagarna 2005, Lund University, May 13-14, 2005

NNLO (Two Loops) in the Chiral Expansion

- Collaboration with Johan Bijnens and Niclas Danielsson at Lund University;
- Calculation of pseudoscalar meson masses and decay constants to NNLO, or $\mathcal{O}(p^6)$ in Partially Quenched Chiral Perturbation Theory (PQ χ PT). Electromagnetic form factors are also under consideration.
- Numerical implementation for the benefit of the Lattice QCD community; Chiral extrapolations to physical sea quark masses.
- Application to upcoming numerical Lattice QCD simulations with light dynamical sea quarks, ultimate objective: → Determination of QCD low-energy constants!

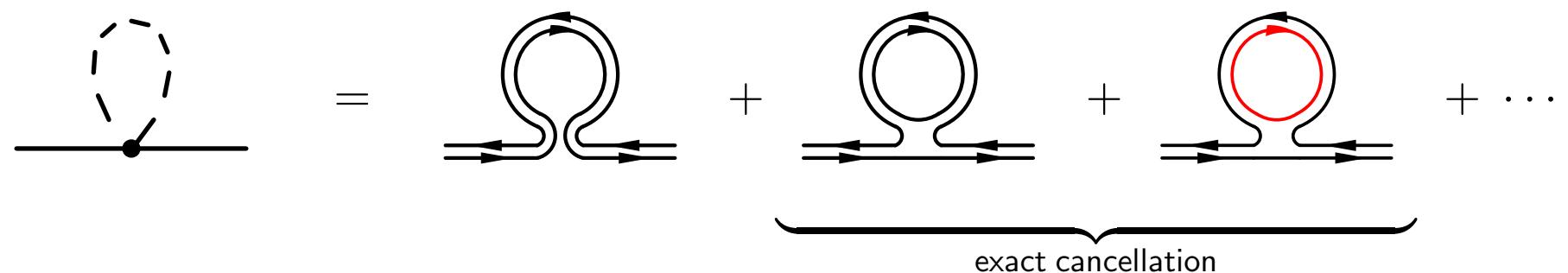
Quenching of Sea Quark Loops in QCD

- Valence and *sea* quark contributions are *treated separately* in Lattice QCD - no such distinction in χ PT.



- Sea quark loops (dynamical fermions) are computationally very demanding; → *Quenched* (no sea quarks) and *partially quenched* (heavy sea quarks) simulations attempted.
- Lattice QCD simulations with *physical u, d* sea quark masses are not (yet?) possible → *Partially Quenched χ PT* applicable if sea quarks are light enough!

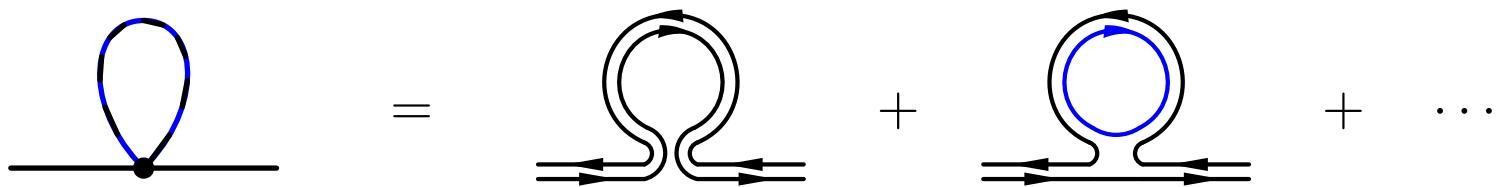
- A first step toward Partially Quenched theories is to remove the contributions with disconnected quark lines from the meson loops:



- Manual intervention at the level of each Feynman diagram is a possible approach; → **Systematic treatment preferable!**
- Introduce **bosonic ghost quarks** in order to cancel out all sea quark contributions from χ PT; → **Supersymmetric formulation¹** of quenched χ PT.

¹C.W. Bernard and M.F.L. Golterman, Phys.Rev. **D49** (1994) 486

- Addition of explicit sea quark fields with independent, arbitrary masses to quenched χ PT; \longrightarrow Partially Quenched χ PT.



- $\text{PQ}\chi\text{PT}$ is a systematic Lagrangian framework² where all necessary cancellations are carried out automatically, for any given set of Feynman diagrams.
- Low-energy constants of $\text{PQ}\chi\text{PT}$ are linear combinations of the unquenched ones; \longrightarrow Direct information about the physical world from partially quenched Lattice QCD data!

²S. Sharpe and N. Shores, Phys.Rev. **D64** (2001) 114510, Phys.Rev. **D62** (2000) 094503

Calculation of Meson Masses to NNLO

- The physical squared masses $M_{\text{phys}}^2 = M^2$ are calculated from the pole position of the resummed propagator,

$$M^2 = M_0^2 + \Sigma(M^2).$$

- Up to NNLO, in terms of the lowest order mass M_0 and the self-energy contribution Σ , the masses of the pseudoscalar mesons are given by

$$M_{\text{phys}}^2 = M_0^2 + \Sigma_4(M_0^2) + \underbrace{\Sigma_4(M_0^2) \frac{\partial \Sigma_4(p^2)}{\partial p^2} \Big|_{M_0^2}}_{\mathcal{O}(p^6) \text{ contribution}} + \Sigma_6(M_0^2) + \mathcal{O}(p^8).$$

- The self-energy Σ_4 represents the NLO (one-loop) mass shift.

Calculation of Decay Constants to NNLO

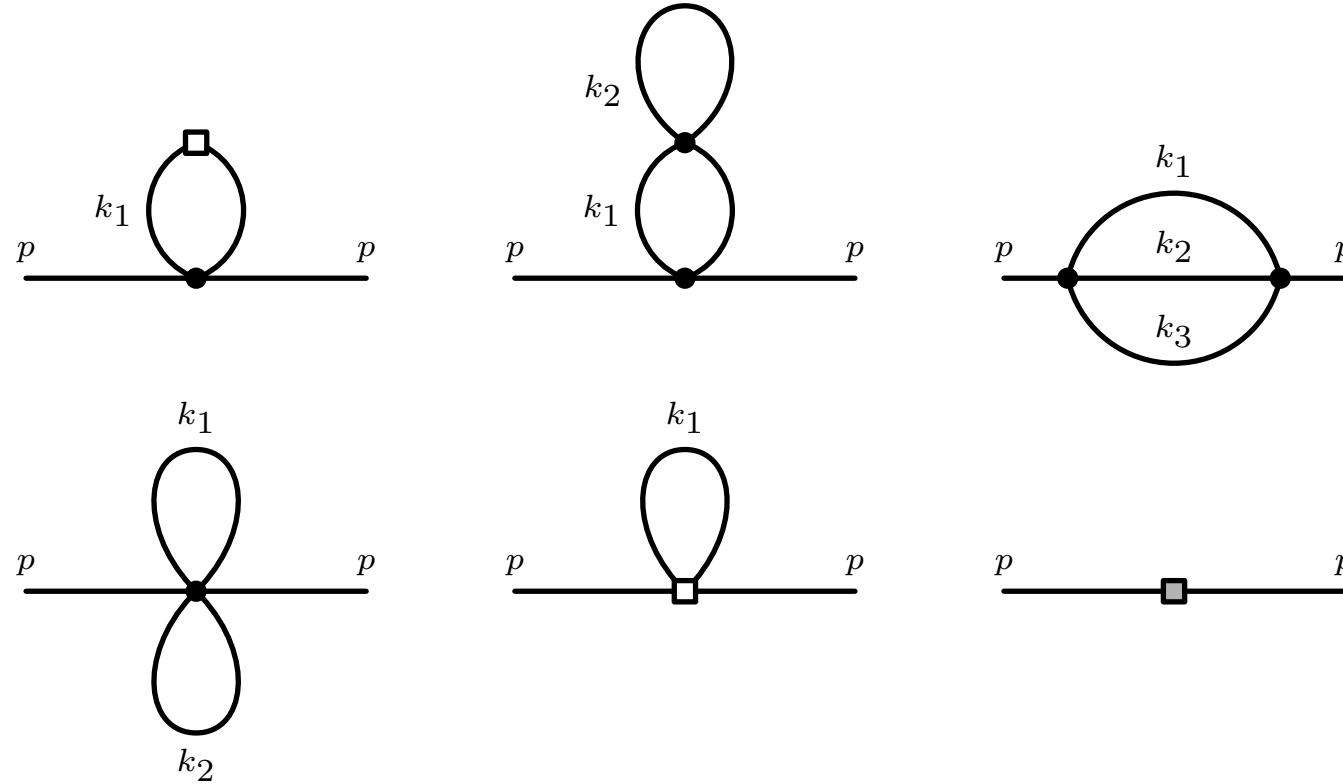
- The decay constants F^a are calculated from the matrix element of the axial current operator,

$$\langle 0 | A_\mu^a(0) | \phi^a(p) \rangle = i\sqrt{2} p_\mu F^a,$$

- Up to NNLO, in terms of the lowest order result $F_2 = F_0$ and the self-energy contributions Σ from the wavefunction renormalization, one obtains

$$\begin{aligned} F_{\text{phys}} &= \underbrace{F_0 + F_4(M_0^2) + F_0 \frac{\partial \Sigma_4(p^2)}{2 \partial p^2} \Big|_{M_0^2}}_{\mathcal{O}(p^4) \text{ contribution}} + F_0 \frac{3}{8} \left(\frac{\partial \Sigma_4(p^2)}{\partial p^2} \Big|_{M_0^2} \right)^2 \\ &\quad + F_0 \frac{\partial \Sigma_6(p^2)}{2 \partial p^2} \Big|_{M_0^2} + F_4(M_0^2) \frac{\partial \Sigma_4(p^2)}{2 \partial p^2} \Big|_{M_0^2} + F_6(M_0^2) + \mathcal{O}(p^8). \end{aligned}$$

- At NNLO, the following diagrams contribute to $-i\Sigma$:



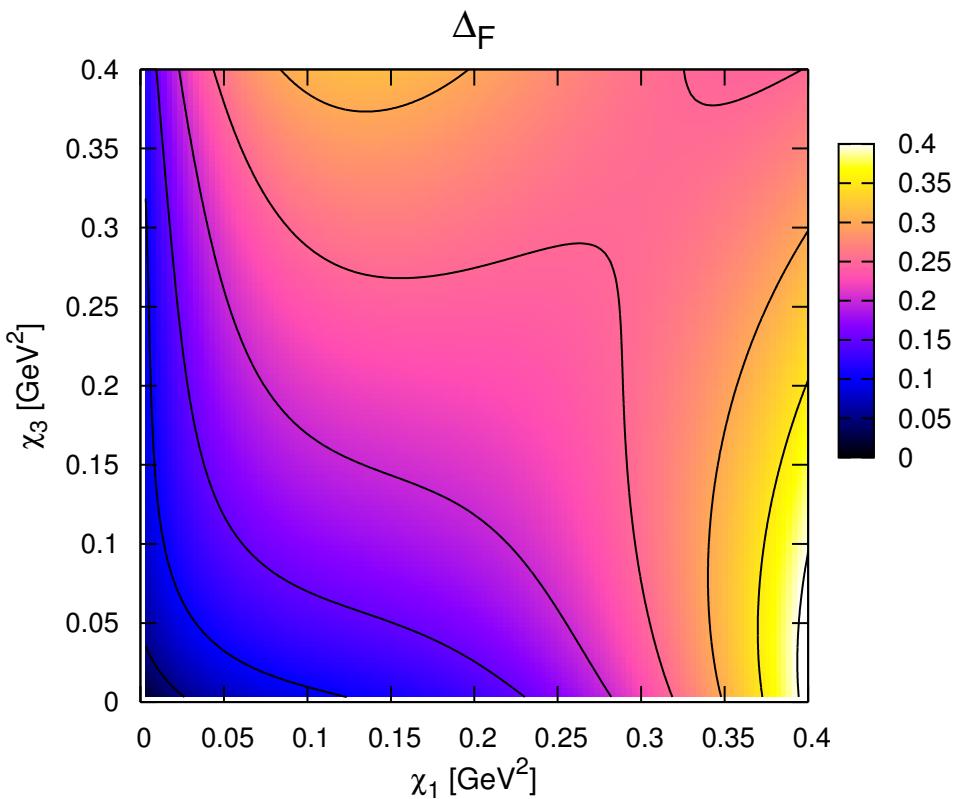
- Black dots, open squares and shaded squares denote vertices of the $\mathcal{O}(p^2)$, $\mathcal{O}(p^4)$ and $\mathcal{O}(p^6)$ Lagrangians, respectively.

Technical Issues and Numerical Results

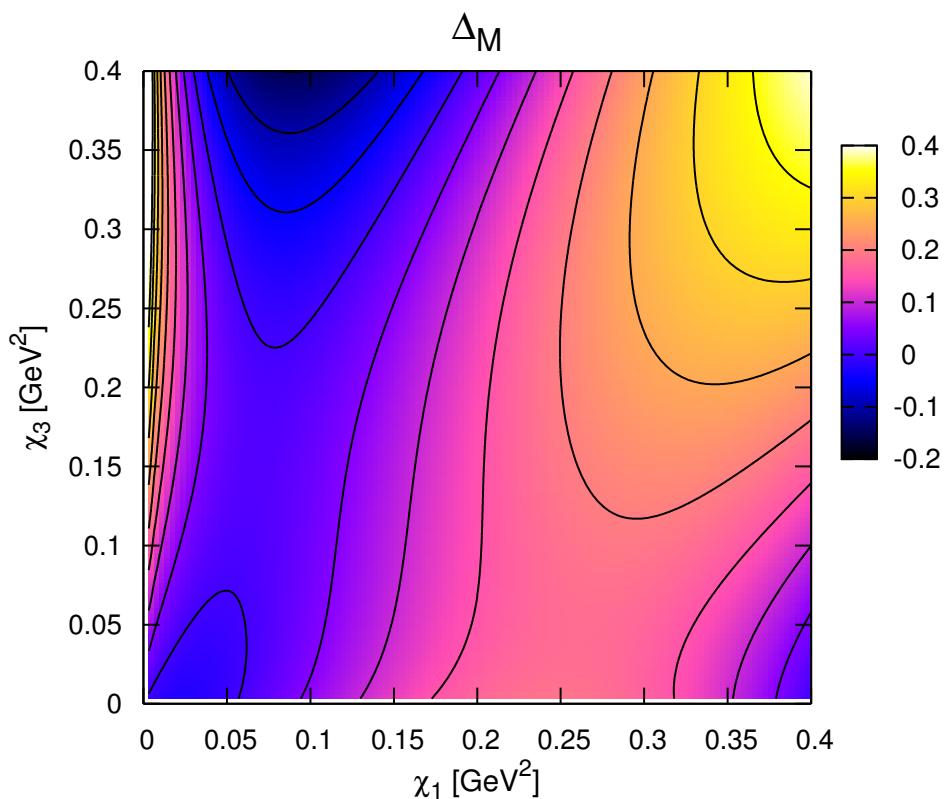
- Expressions in NNLO PQ χ PT are very large; → Heavy use of FORM³.
- The raw FORM output is ugly, bloated and inefficient and contains $\sim 10^4$ terms; → Efficient simplification with MAPLE to about 350 terms!
- The NNLO expressions for F and M^2 depend on several low-energy constants: F_0, B_0 of $\mathcal{O}(p^2)$, the L_i of $\mathcal{O}(p^4)$, and the K_i of $\mathcal{O}(p^6)$.
- Quark mass dependence through "squared masses" $\chi_i = 2 B_0 m_i$.
- Numerical results given in terms of the relative corrections Δ_i , defined as $M^2 = M_0^2 (1 + \Delta_M)$ and $F = F_0 (1 + \Delta_F)$.

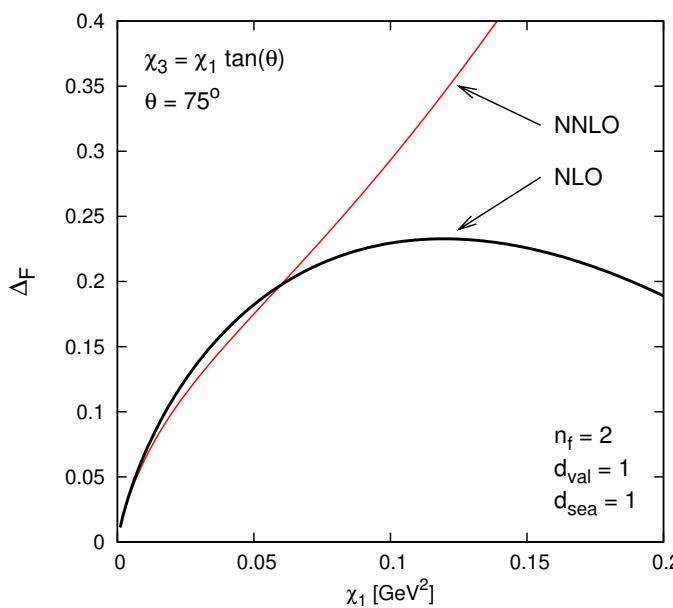
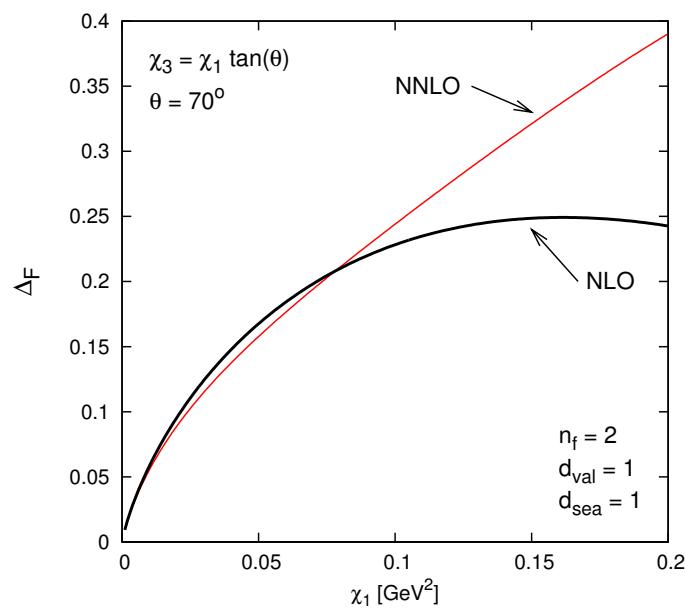
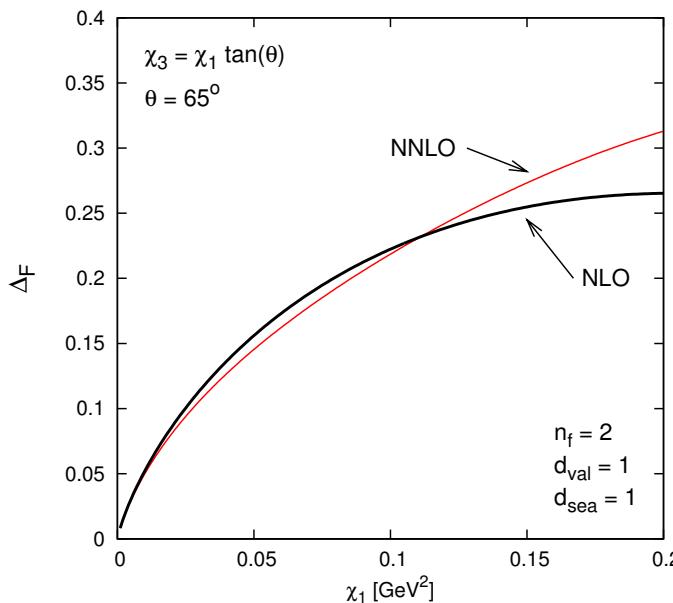
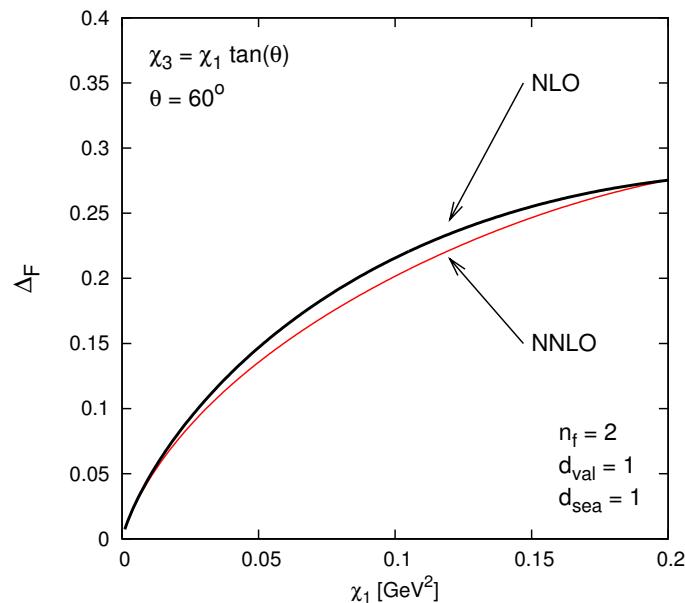
³J. Vermaseren, <http://www.nikhef.nl/~form>

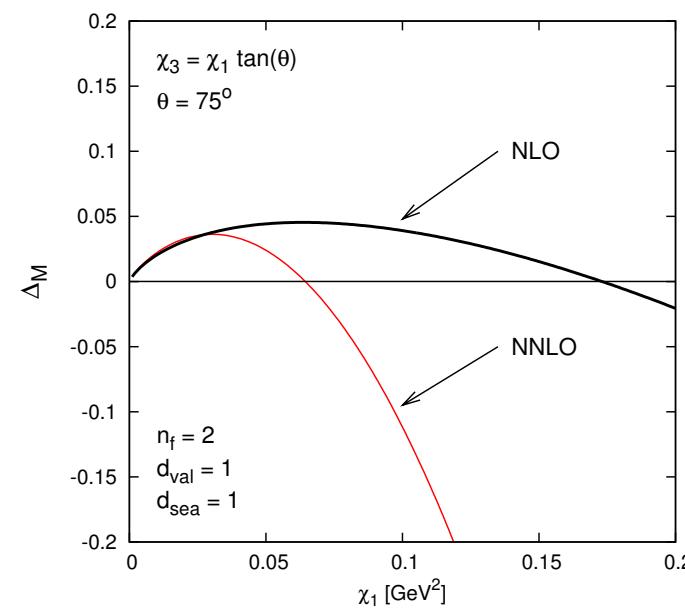
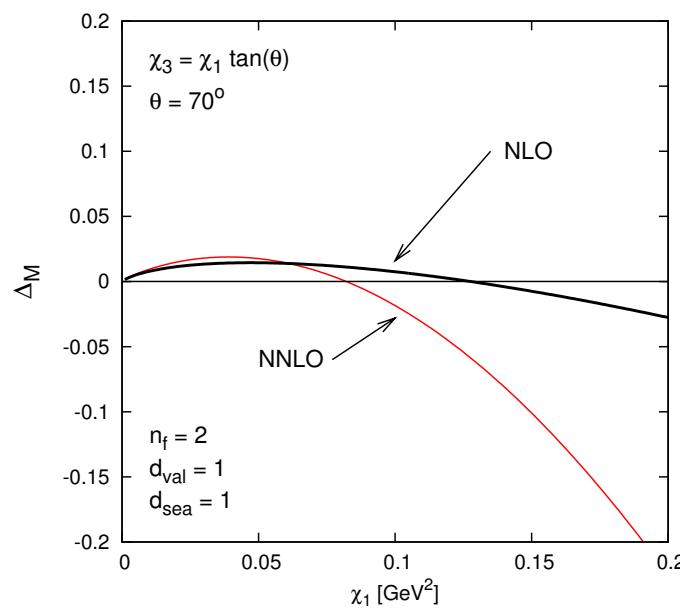
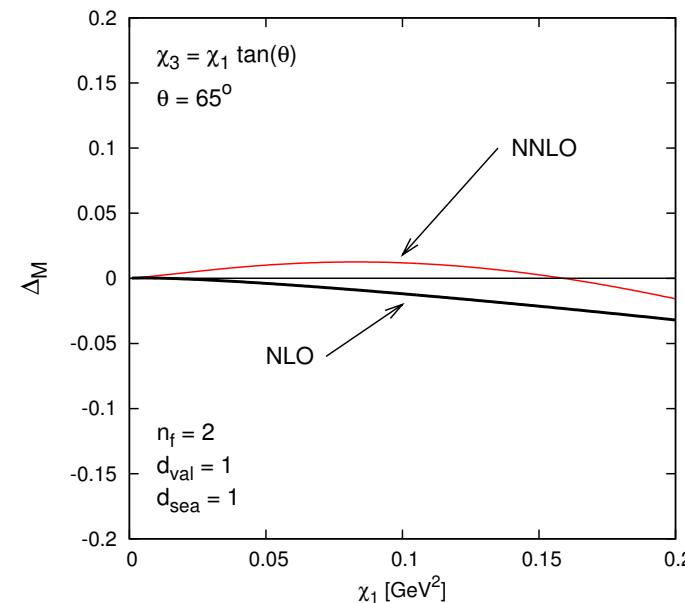
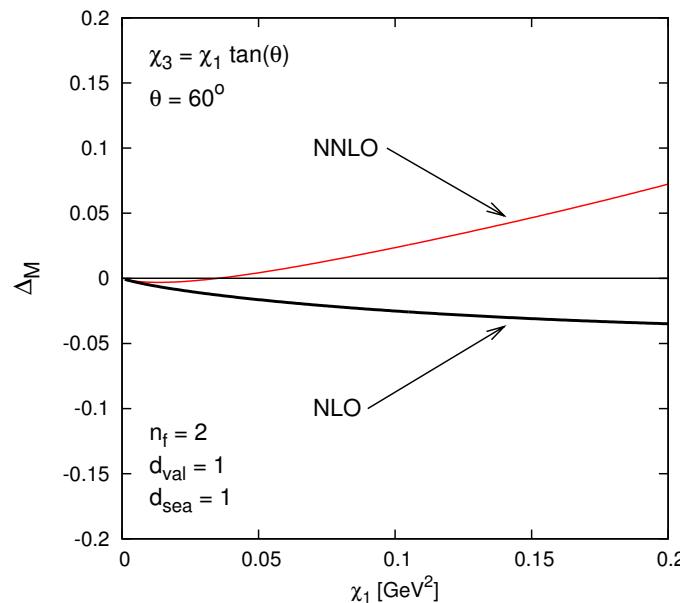
- NNLO contribution to the decay constant Δ_F , defined as $F = F_0(1 + \Delta_F)$:
- Calculated for two sea quark flavors, $n_{\text{sea}} = 2$.
- Plotted for degenerate sea quark masses χ_3 .
- Unquenched χ PT recovered along the diagonal!



- NNLO contribution to the meson mass Δ_M , defined as $M^2 = M_0^2 (1 + \Delta_M)$:
- Note the divergence of the quenched chiral logarithms as $\chi_1 \rightarrow 0$ for a constant sea quark mass.
- Calculated for two sea quark flavors, $n_{\text{sea}} = 2$.
- Plotted for degenerate sea quark masses χ_3 .







Published Papers in PQ χ PT:

- Phys.Rev.**D70**:111503 (2004), hep-lat/0406017,
"Pseudoscalar meson mass to two loops in three-flavor partially quenched chiral perturbation theory"
- Phys.Rev.**D71**:094502 (2005), hep-lat/0501014,
"Decay constants of pseudoscalar mesons to two loops in three-flavor partially quenched chiral perturbation theory"

To appear shortly:

- Meson masses and decay constants in two-flavor PQ χ PT,
- Meson masses in three-flavor PQ χ PT, for nondegenerate sea quarks.

Planned future calculations (depending on interest):

- Electromagnetic form factors in PQ χ PT, properties of neutral pseudoscalars.