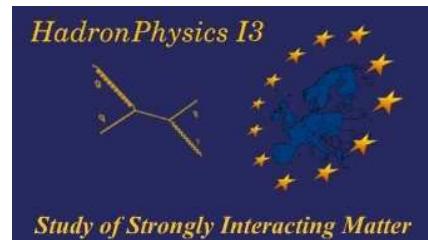




LUND
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RECENT ISSUES IN ETAS AND KAONS

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Overview

- Motivation: Standard Model
- Method: Chiral Perturbation Theory
- Two examples:
 - $\eta \rightarrow 3\pi$: $m_u - m_d$
 - $K_{\ell 3}$: V_{us}
- Conclusions

The Standard Model

The Standard Model Lagrangian has four parts:

$$\underbrace{\mathcal{L}_H(\phi)}_{\text{Higgs}} + \underbrace{\mathcal{L}_G(W, Z, G)}_{\text{Gauge}}$$

$$\underbrace{\sum_{\psi=\text{fermions}} \bar{\psi} i \not{D} \psi}_{\text{gauge-fermion}} + \underbrace{\sum_{\psi, \psi'=\text{fermions}} g_{\psi\psi'} \bar{\psi} \phi \psi'}_{\text{Yukawa}}$$

The Standard Model

What is tested ?

gauge-fermion	Very well tested
Higgs	Limits only, real tests coming up
Gauge	Well tested, QCD at low-energy nonperturbative
Yukawa	Flavour Physics

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Discrete symmetries:

- C Charge Conjugation
- P Parity
- T Time Reversal

QCD and QED conserve C,P,T separately,
Weak breaks C and P, only Yukawa breaks CP
Field theory implies CPT

Standard Model

We want to study Low-energy QCD as well as possible and determine couplings in the Yukawa sector as precisely as possible

Effective Field Theory

Main Ideas:

- Use right degrees of freedom : essence of (most) physics
- If mass-gap in the excitation spectrum: neglect degrees of freedom above the gap.

Examples:

Solid state physics: conductors: neglect the empty bands above the partially filled one
Atomic physics: Blue sky: neglect atomic structure

Power Counting

- ➡ gap in the spectrum \Rightarrow separation of scales
- ➡ with the lower degrees of freedom, build the most general effective Lagrangian

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- Need some ordering principle: power counting

Power Counting

- gap in the spectrum \Rightarrow separation of scales
 - with the lower degrees of freedom, build the most general effective Lagrangian
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- \rightarrow $\infty \#$ parameters
 - \rightarrow Where did my predictivity go ?
-
- 
- Need some ordering principle: power counting
-
- \rightarrow Taylor series expansion does not work (convergence radius is zero)
 - \rightarrow Continuum of excitation states need to be taken into account

Why Field Theory ?

- ➡ Only known way to combine QM and special relativity
- ➡ Off-shell effects: there as new free parameters

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Drawbacks

- Many parameters (but finite number at any order)
any model has few parameters but model-space is large
- expansion: it might not converge or only badly

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Advantages

- Calculations are (relatively) simple
- It is general: model-independent
- **Theory** \implies errors can be estimated
- Systematic: ALL effects at a given order can be included
- Even if no convergence: classification of models often useful

Chiral Perturbation Theory

Chiral Symmetry

QCD: 3 light quarks: equal mass: interchange: $SU(3)_V$

But $\mathcal{L}_{QCD} = \sum_{q=u,d,s} [i\bar{q}_L \not{D} q_L + i\bar{q}_R \not{D} q_R - m_q (\bar{q}_R q_L + \bar{q}_L q_R)]$

So if $m_q = 0$ then $SU(3)_L \times SU(3)_R$.

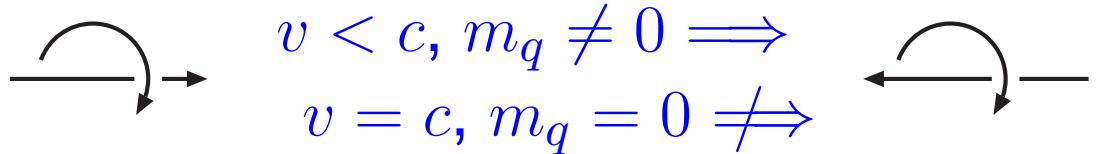
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Can also see that via 

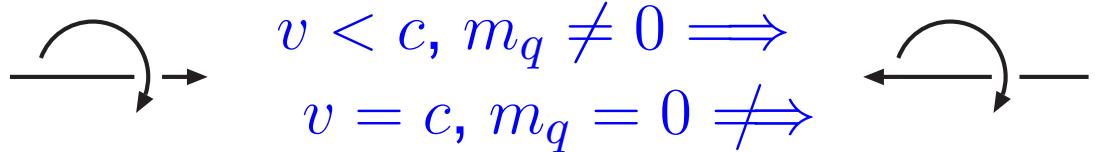
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So if $m_q = 0$ then $SU(3)_L \times SU(3)_R$.

Can also see that via  $v < c, m_q \neq 0 \Rightarrow$
 $v = c, m_q = 0 \Rightarrow$

Degrees of freedom: Goldstone Bosons from Chiral Symmetry Spontaneous Breakdown

Power counting: Dimensional counting

Expected breakdown scale: Resonances, so M_ρ or higher depending on the channel

Chiral Perturbation Theory

$$\langle \bar{q}q \rangle = \langle \bar{q}_L q_R + \bar{q}_R q_L \rangle \neq 0$$

$SU(3)_L \times SU(3)_R$ broken spontaneously to $SU(3)_V$

8 generators broken \implies 8 massless degrees of freedom
and interaction vanishes at zero momentum

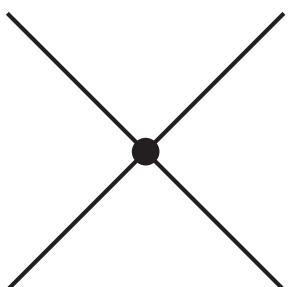
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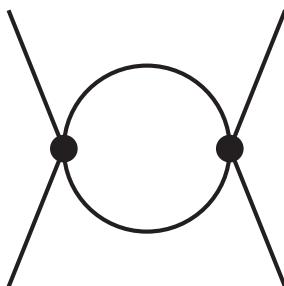
8 generators broken \implies 8 massless degrees of freedom
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Power counting in momenta:



$$p^2$$

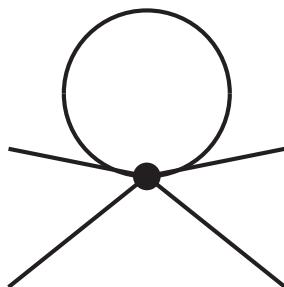
$$\int d^4 p$$



$$1/p^2$$

$$p^4$$

$$(p^2)^2 (1/p^2)^2 p^4 = p^4$$



$$(p^2) (1/p^2) p^4 = p^4$$

Chiral Perturbation Theory

Perturbation Theory with mesons, expansion in momenta

The physics of pions, kaons and eta

Actually useful:

- tree level: sixties
- one-loop: eighties
- two-loop now (and especially in Lund ☺)
- Similar techniques: technicolour, little Higgs,...

Remember: Eta Physics Handbook

Physica Scripta, Vol. T99, 2002

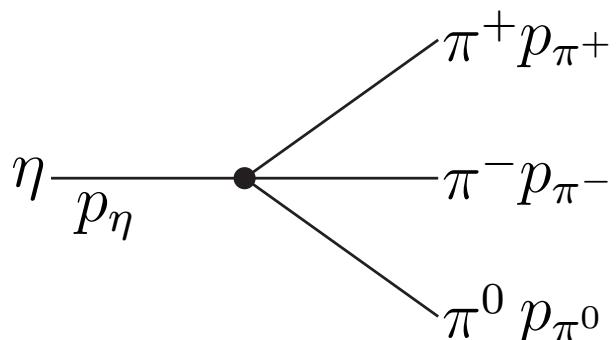
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$\eta \rightarrow 3\pi$ beyond p^4 : Basic

Review: JB, Gasser, Phys.Scripta T99(2002)34 [hep-ph/0202242]



$$\begin{aligned} s &= (p_{\pi^+} + p_{\pi^-})^2 = (p_\eta - p_{\pi^0})^2 \\ t &= (p_{\pi^-} + p_{\pi^0})^2 = (p_\eta - p_{\pi^+})^2 \\ u &= (p_{\pi^+} + p_{\pi^0})^2 = (p_\eta - p_{\pi^-})^2 \end{aligned}$$

$$s + t + u = m_\eta^2 + 2m_{\pi^+}^2 + m_{\pi^0}^2 \equiv 3s_0.$$

$$\langle \pi^0 \pi^+ \pi^- \text{out} | \eta \rangle = i (2\pi)^4 \delta^4 (p_\eta - p_{\pi^+} - p_{\pi^-} - p_{\pi^0}) A(s, t, u).$$

$$\langle \pi^0 \pi^0 \pi^0 \text{out} | \eta \rangle = i (2\pi)^4 \delta^4 (p_\eta - p_1 - p_2 - p_3) \overline{A}(s_1, s_2, s_3)$$

$$\overline{A}(s_1, s_2, s_3) = A(s_1, s_2, s_3) + A(s_2, s_3, s_1) + A(s_3, s_1, s_2),$$

$\eta \rightarrow 3\pi$ beyond p^4 : Lowest order

Pions are in $I = 1$ state $\Rightarrow A \sim (m_u - m_d)$ or α_{em}

- α_{em} effect is small (but large via $m_{\pi^+} - m_{\pi^0}$)
- $\eta \rightarrow \pi^+ \pi^- \pi^0 \gamma$ needs to be included directly

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ChPT:Cronin 67: $A(s, t, u) = \frac{B_0(m_u - m_d)}{3\sqrt{3}F_\pi^2} \left\{ 1 + \frac{3(s - s_0)}{m_\eta^2 - m_\pi^2} \right\}$

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or with $Q^2 \equiv \frac{m_s^2 - \hat{m}^2}{m_d^2 - m_u^2}$, $\hat{m} = \frac{1}{2}(m_u + m_d)$

$$A(s, t, u) = \frac{1}{Q^2} \frac{m_K^2}{m_\pi^2} (m_\pi^2 - m_K^2) \frac{M(s, t, u)}{3\sqrt{3}F_\pi^2},$$

with at lowest order $M(s, t, u) = \frac{3s - 4m_\pi^2}{m_\eta^2 - m_\pi^2}$.

$\eta \rightarrow 3\pi$ beyond p^4 : p^2 and p^4

$\Gamma(\eta \rightarrow 3\pi) \propto |A|^2 \propto Q^{-4}$ allows a PRECISE measurement

$Q \approx 24$ gives lowest order $\Gamma(\eta \rightarrow \pi^+ \pi^- \pi^0) \approx 66 \text{ eV}$.

Other Source from $m_{K^+}^2 - m_{K^0}^2 \sim Q^{-2} \Rightarrow Q = 20.0 \pm 1.5$

Lowest order prediction $\Gamma(\eta \rightarrow \pi^+ \pi^- \pi^0) \approx 140 \text{ eV}$.

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Lowest order prediction $\Gamma(\eta \rightarrow \pi^+ \pi^- \pi^0) \approx 140 \text{ eV}$.

At order p^4 :
$$\frac{\int dLIPS |A_2 + A_4|^2}{\int dLIPS |A_2|^2} = 2.4,$$

(LIPS=Lorentz invariant phase-space)

Major source: large S -wave final state rescattering

Experiment: $\Gamma(\eta \rightarrow 3\pi) = 295 \pm 17 \text{ eV}$.

$\eta \rightarrow 3\pi$ beyond p^4 : Dispersive

Try to resum the S -wave rescattering:

Anisovich-Leutwyler (AL), Kambor,Wiesendanger,Wyler (KWW)

Different method but similar approximations

Here: simplified version of AL

Up to p^8 : No absorptive parts from $\ell \geq 2$

$$\Rightarrow M(s, t, u) =$$

$$M_0(s) + (s - u)M_1(t) + (s - t)M_1(t) + M_2(t) + M_2(u) - \frac{2}{3}M_2(s)$$

M_I : “roughly” contributions with isospin 0,1,2

$\eta \rightarrow 3\pi$ beyond p^4 : Dispersive

3 body dispersive: difficult: keep only 2 body cuts

start from $\pi\eta \rightarrow \pi\pi$ ($m_\eta^2 < 3m_\pi^2$) standard dispersive analysis
analytically continue to physical m_η^2 .

$$M_I(s) = \frac{1}{\pi} \int_{4m_\pi^2}^{\infty} ds' \frac{\text{Im}M_I(s')}{s' - s - i\varepsilon}$$

$$\text{Im}M_I(s') \longrightarrow \text{disc}M_I(s) = \frac{1}{2i} (M_I(s + i\varepsilon) - M_I(s - i\varepsilon))$$

$$M_0(s) = a_0 + b_0s + c_0s^2 + \frac{s^3}{\pi} \int \frac{ds'}{s'^3} \frac{\text{disc}M_0(s')}{s' - s - i\varepsilon},$$

$$M_1(s) = a_1 + b_1s + \frac{s^2}{\pi} \int \frac{ds'}{s'^2} \frac{\text{disc}M_1(s')}{s' - s - i\varepsilon},$$

$$M_2(s) = a_2 + b_2s + c_2s^2 + \frac{s^3}{\pi} \int \frac{ds'}{s'^3} \frac{\text{disc}M_2(s')}{s' - s - i\varepsilon}.$$

$\eta \rightarrow 3\pi$ beyond p^4

AL: Lowest order is $M(s, t, u) = \frac{3s - 4m_\pi^2}{m_\eta^2 - m_\pi^2}$

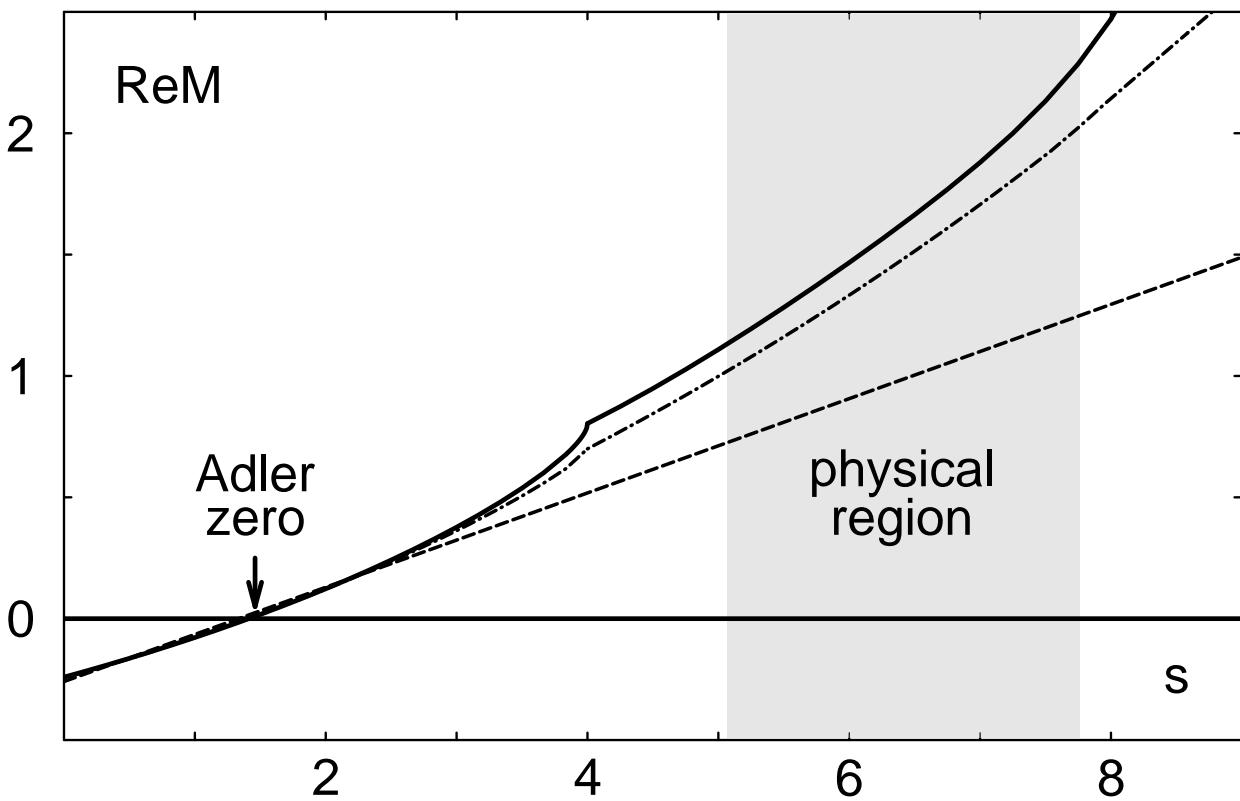
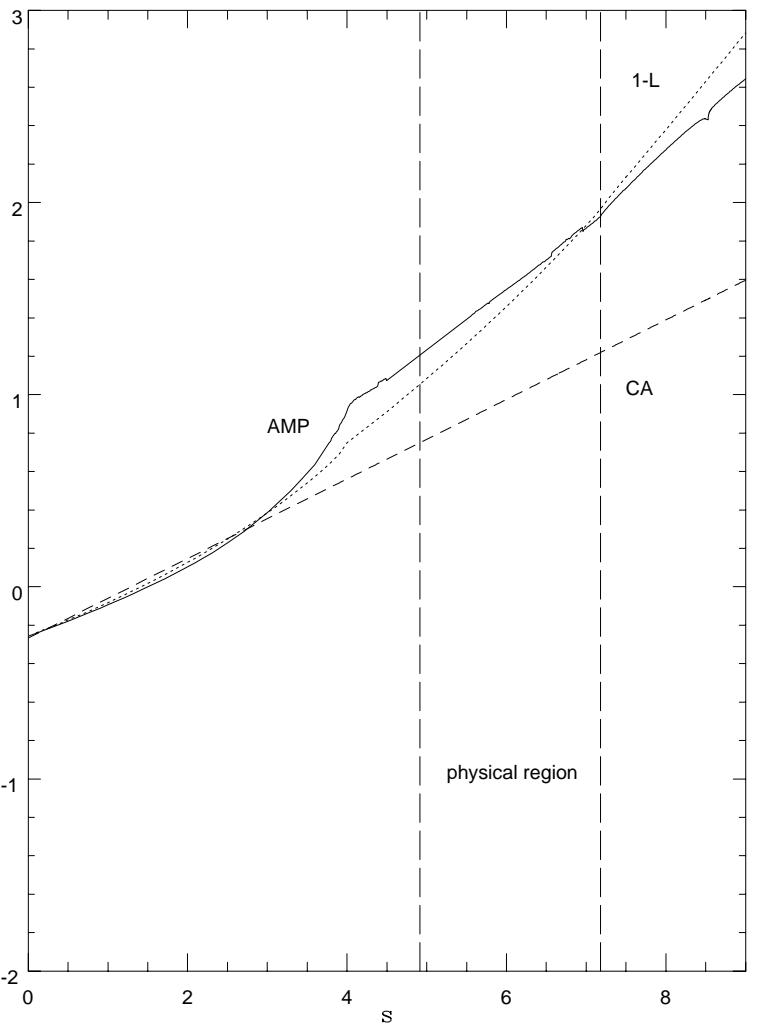
zero at $s_4/3 m_\pi^2$: remains in the neighbourhood:
match position of s_A and slope of Adler zero.

KWW: fix amplitude at some place(s) in s, t, u plane to be equal to p^4

Both find moderate increases over p^4 : $\sim 15\%$ in amplitude

Dalitzplot distributions provide a check

$\eta \rightarrow 3\pi$ beyond p^4



Two Loop Calculation: why

- In $K_{\ell 4}$ dispersive gave about half of p^6 in amplitude
- Same order in ChPT as masses for consistency check on m_u/m_d
- Check size of 3 pion dispersive part
- Technology exists:
 - Two-loops: Amorós,JB,Dhonte,Talavera,...
 - Dealing with the mixing π^0 - η :
Amorós,JB,Dhonte,Talavera 01

Status

$$A(s, t, u) = \sin \epsilon M(s, t, u)$$

$$\epsilon \approx \frac{\sqrt{3}}{4} \frac{m_d - m_u}{m_s - \hat{m}}$$

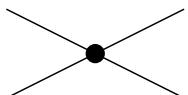
$$M(s, t, u) =$$

$$M_0(s) + (s - u)M_1(t) + (s - t)M_1(t) + M_2(t) + M_2(u) - \frac{2}{3}M_2(s)$$

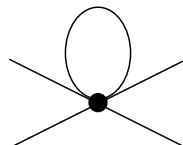
$$M^{(2)} = \frac{1}{F_\pi^2} \left(\frac{4}{3} m_\pi^2 - s \right)$$

	JB	Ghorbani	Talavera
p^4	All agree and with GL		
$p^6 \infty$	All cancel		
p^6 Vertex	Checked and agree	in progress	
p^6 Rest	Comparison in progress		
p^6 Numerics	Very preliminary		Very preliminary

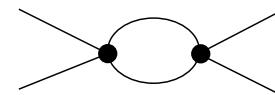
Diagrams



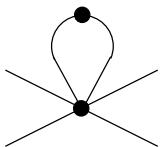
(a)



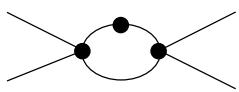
(b)



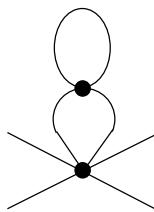
(c)



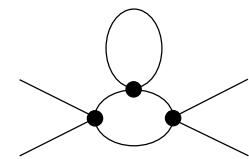
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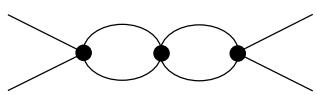
(e)



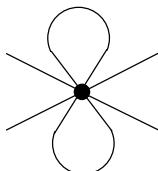
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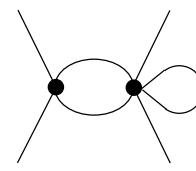
(g)



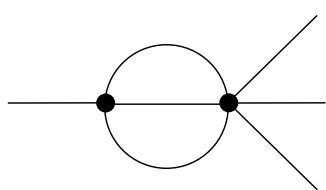
(h)



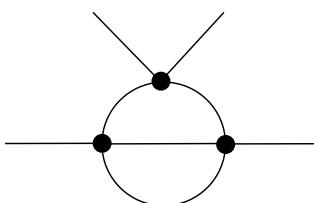
(i)



(j)

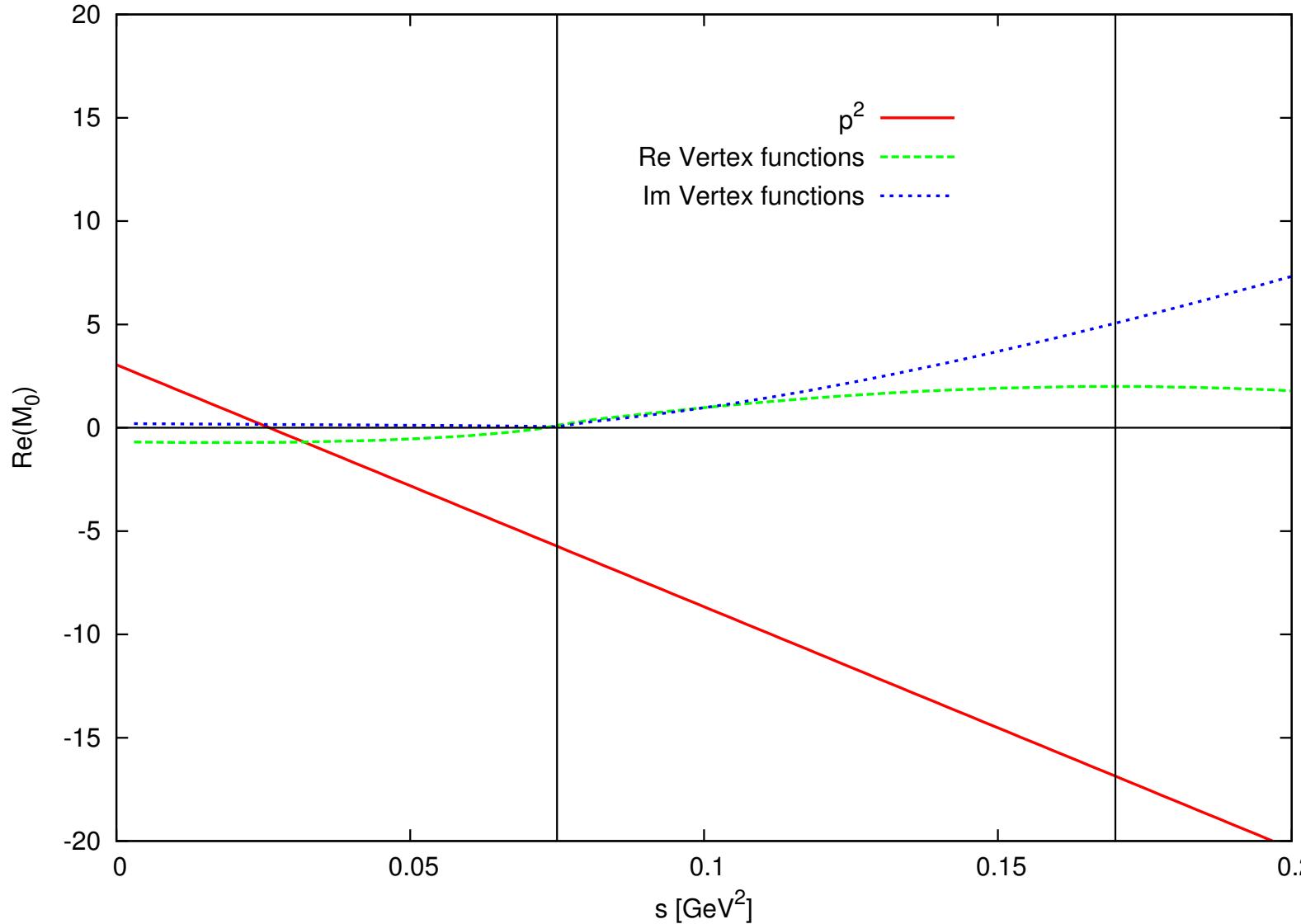


(h)

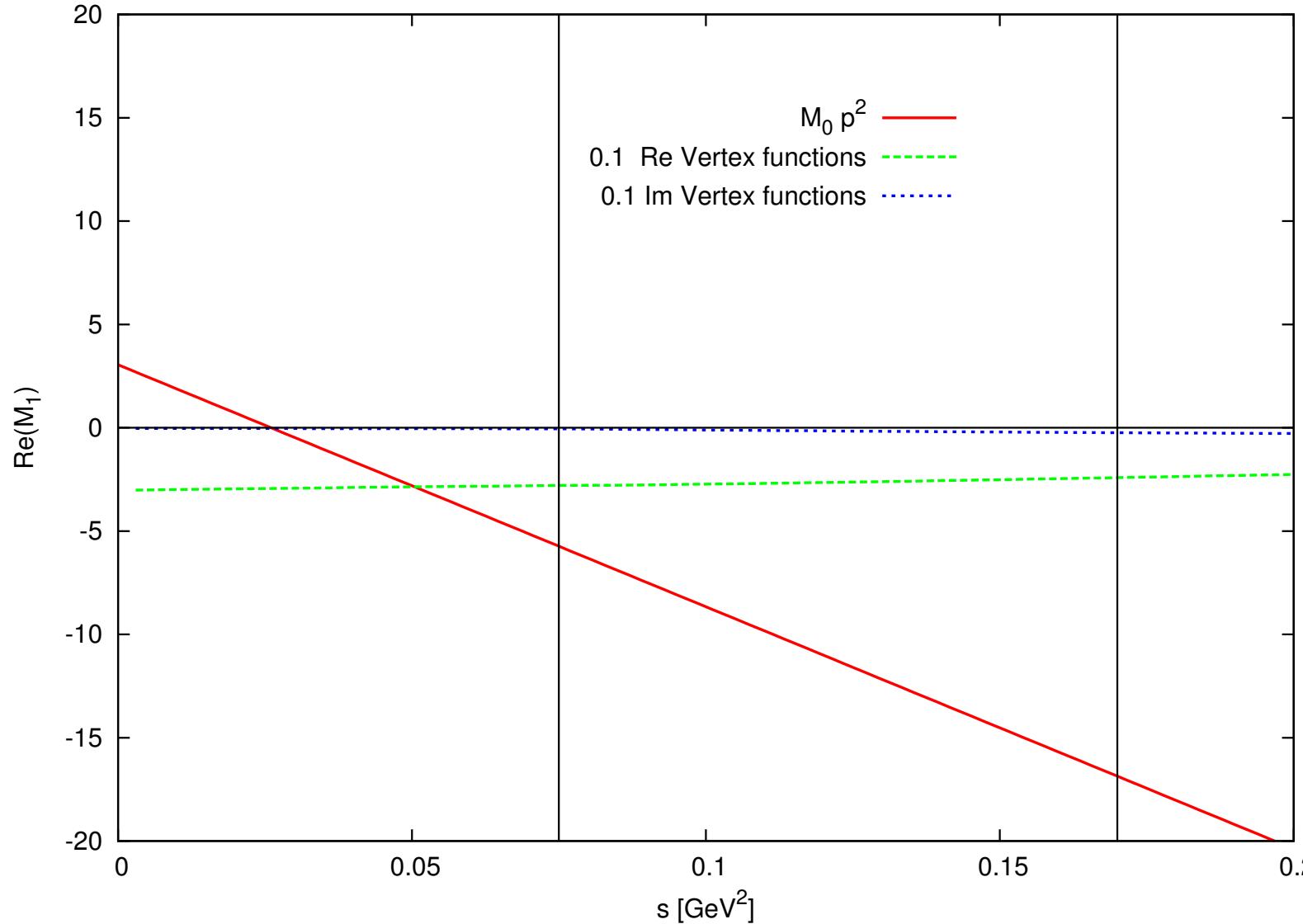


(i)

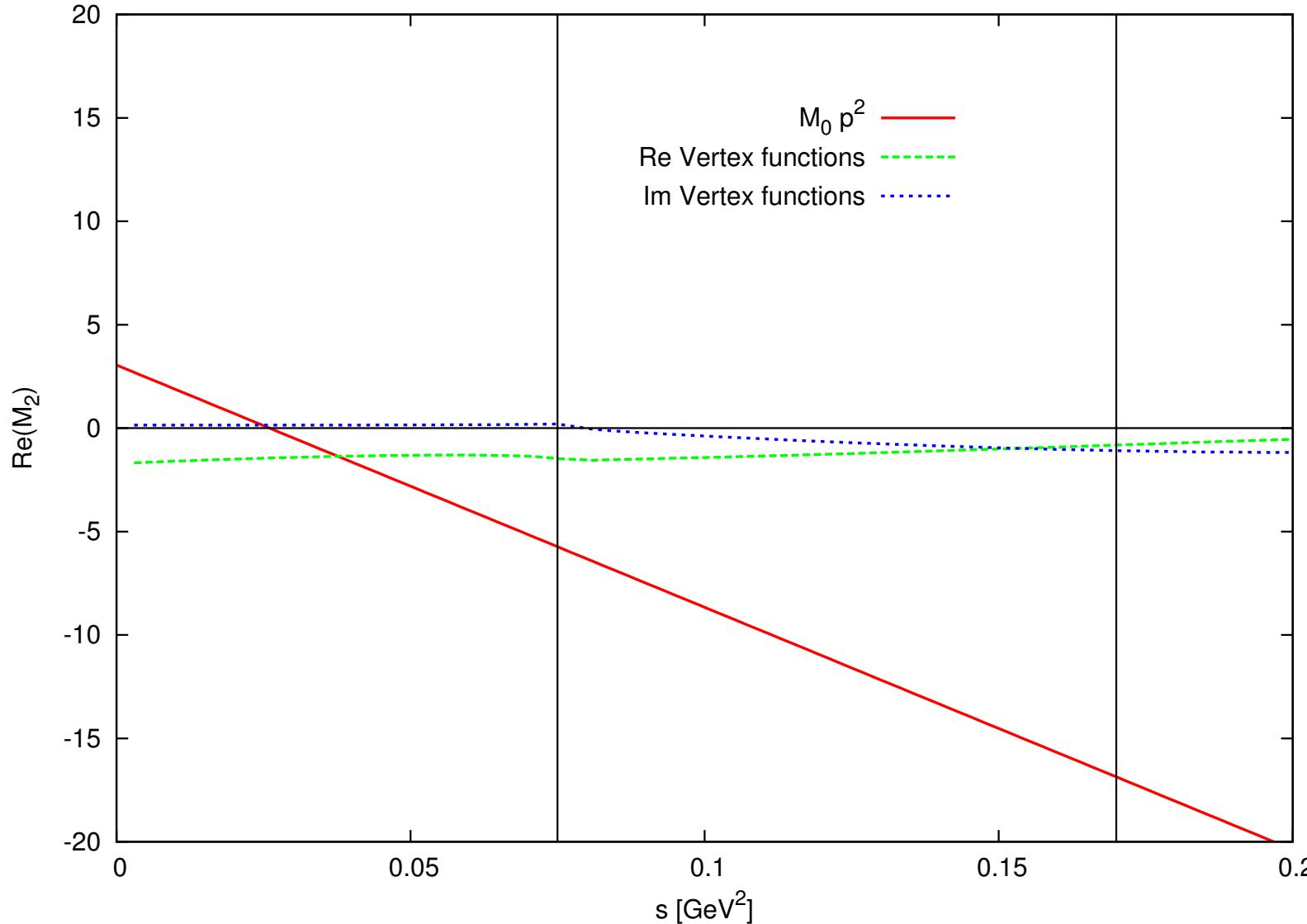
M_0 : Vertex function contribution



M_1 : Vertex function contribution



M_2 : Vertex function contribution



$K_{\ell 3}$ Definitions

$$K_{\ell 3}^+ : \quad K^+(p) \rightarrow \pi^0(p') \ell^+(p_\ell) \nu_\ell(p_\nu)$$

$$K_{\ell 3}^0 : \quad K^0(p) \rightarrow \pi^-(p') \ell^+(p_\ell) \nu_\ell(p_\nu)$$

$$K_{\ell 3}^+ : \quad T = \frac{G_F}{\sqrt{2}} V_{us}^\star \ell^\mu F_\mu^+(p', p)$$

$$\ell^\mu = \bar{u}(p_\nu) \gamma^\mu (1 - \gamma_5) v(p_\ell)$$

$$\begin{aligned} F_\mu^+(p', p) &= <\pi^0(p') \mid V_\mu^{4-i5}(0) \mid K^+(p)> \\ &= \frac{1}{\sqrt{2}} [(p' + p)_\mu f_+^{K^+\pi^0}(t) + (p - p')_\mu f_-^{K^+\pi^0}(t)] \end{aligned}$$

$$\text{Isospin: } f_+^{K^0\pi^-}(t) = f_+^{K^+\pi^0}(t) = f_+(t)$$

$$f_-^{K^0\pi^-}(t) = f_-^{K^+\pi^0}(t) = f_-(t)$$

$K_{\ell 3}$ Definitions and V_{us}

Scalar formfactor:

$$f_0(t) = f_+(t) + \frac{t}{m_K^2 - m_\pi^2} f_-(t)$$

Usual parametrization:

$$f_{+,0}(t) = f_+(0) \left(1 + \lambda_{+,0} \frac{t}{m_\pi^2} \right)$$

$|V_{us}|$:

- Know theoretically $f_+(0) = 1 + \dots$
- Short distance correction to G_F from G_μ
Marciano-Sirlin
- Ademollo-Gatto-Behrends-Sirlin theorem:
 $(m_s - \hat{m})^2$
- Isospin Breaking Leuwyler-Roos
- Know experimentally $f_+(0)$

V_{us}

PDG2002:

$$|V_{ud}| = 0.9734 \pm 0.0008 \quad |V_{us}| = 0.2196 \pm 0.0026$$
$$|V_{ud}|^2 + |V_{us}|^2 = (0.9475 \pm 0.0016) + (0.0482 \pm 0.0011) =$$
$$0.9957 \pm 0.0019$$

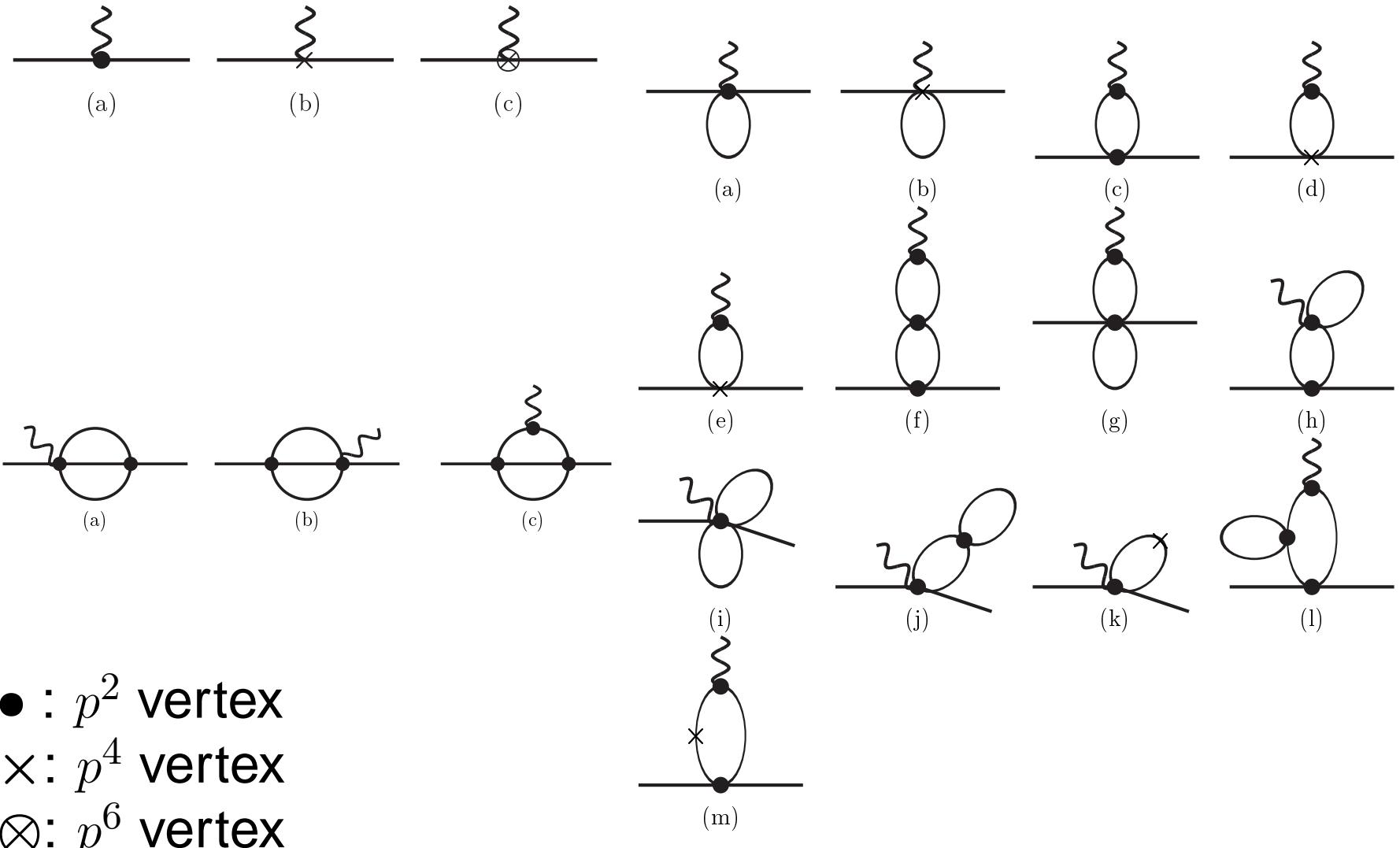
PDG2006:

$$|V_{ud}| = 0.97377 \pm 0.00027 \quad |V_{us}| = 0.2257 \pm 0.0021$$
$$|V_{ud}|^2 + |V_{us}|^2 = (0.94823 \pm 0.00054) + (0.05094 \pm 0.00095) =$$
$$0.99917 \pm 0.00110$$

Problems:

- Ignores $\Delta(0) = 0.0113$ from pure two-loop
- Conflicts between experiments

$K_{\ell 3}$ Diagrams



- : p^2 vertex
- × : p^4 vertex
- ⊗ : p^6 vertex

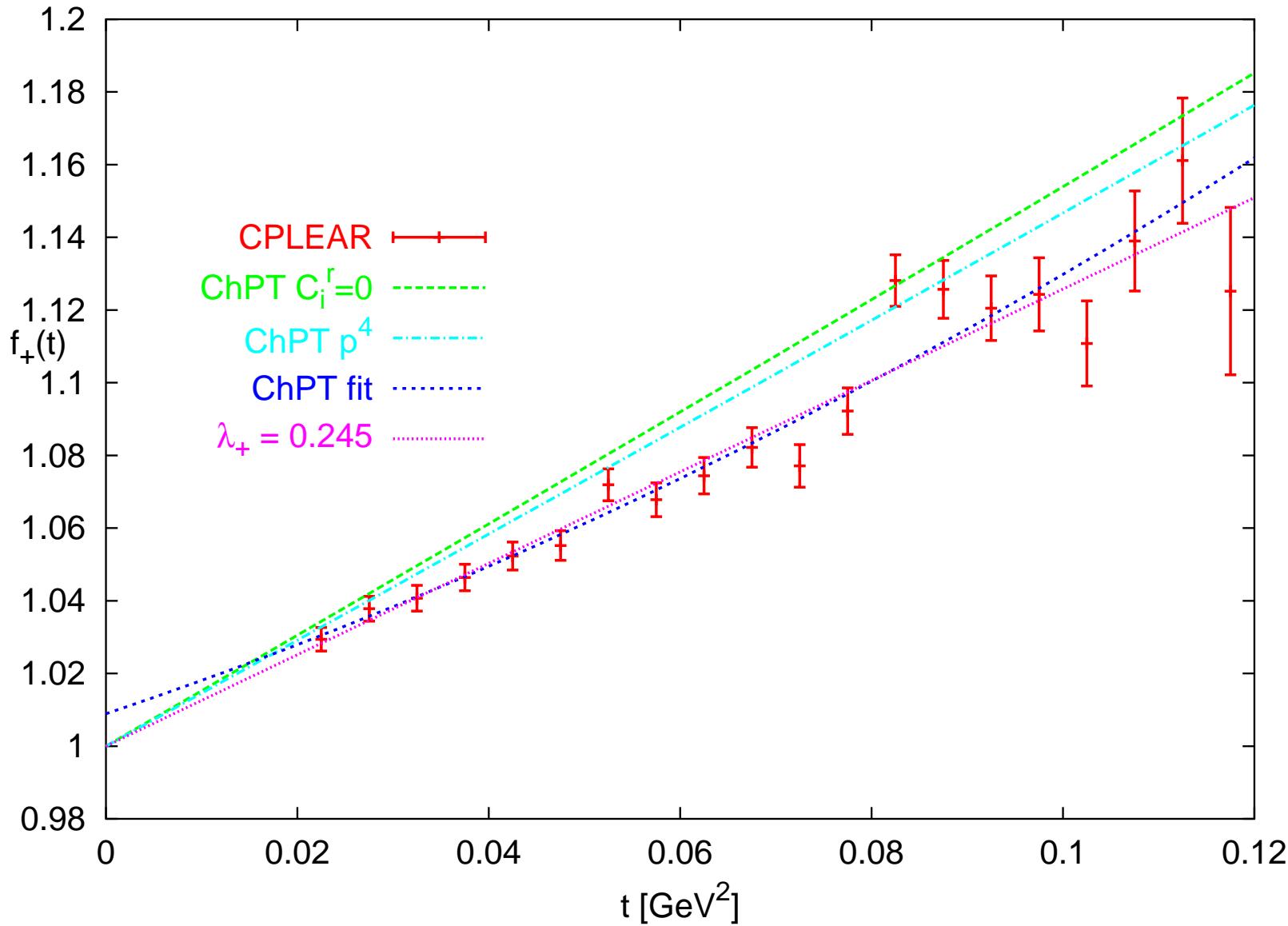
$f_+(t)$ Theory

$$f_+(t) = 1 + f_+^{(4)}(t) + f_+^{(6)}(t)$$

$$f_+^{(4)}(t) = \frac{t}{2F_\pi^2} L_9^r + \text{loops}$$

$$\begin{aligned} f_+^{(6)}(t) = & -\frac{8}{F_\pi^4} (C_{12}^r + C_{34}^r) (m_K^2 - m_\pi^2)^2 + \frac{t}{F_\pi^4} R_{+1}^{K\pi} \\ & + \frac{t^2}{F_\pi^4} (-4C_{88}^r + 4C_{90}^r) + \text{loops}(L_i^r) \end{aligned}$$

ChPT fit to $f_+(t)$

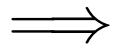


$$f_0(t)$$

Main Result:

$$\begin{aligned}
 f_0(t) = & 1 - \frac{8}{F_\pi^4} (C_{12}^r + C_{34}^r) (m_K^2 - m_\pi^2)^2 \\
 & + 8 \frac{t}{F_\pi^4} (2C_{12}^r + C_{34}^r) (m_K^2 + m_\pi^2) + \frac{t}{m_K^2 - m_\pi^2} (F_K/F_\pi - 1) \\
 & - \frac{8}{F_\pi^4} t^2 C_{12}^r + \bar{\Delta}(t) + \Delta(0).
 \end{aligned}$$

$\bar{\Delta}(t)$ and $\Delta(0)$ contain NO C_i^r and only depend on the L_i^r at order p^6



All needed parameters can be determined experimentally

$$\Delta(0) = -0.0080 \pm 0.0057[\text{loops}] \pm 0.0028[L_i^r].$$

Experiment

Form Factor comparison

KTeV [PRD 70(2004)]

K^0_{e3} quadratic fit: $\lambda''_+ \neq 0$ @ 4σ level

$K^0_{\mu 3}$ quadratic fit: $\lambda_0 = (13.72 \pm 1.31) 10^{-3}$

Slopes consistent for K^0_{e3} and $K^0_{\mu 3}$

ISTRA+ [PLB 581(2004), PLB 589(2004)]

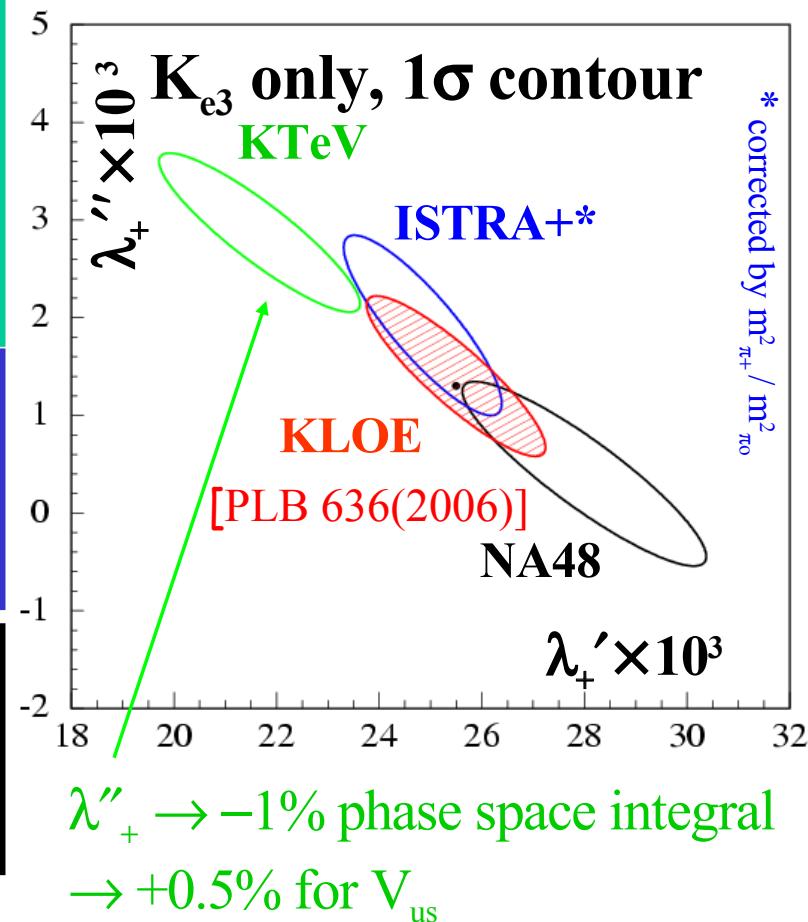
K^-_{e3} quadratic fit: $\lambda''_+ \neq 0$ @ 2σ level

$K^-_{\mu 3}$ quadratic fit: $\lambda_0 = (17.11 \pm 2.31) 10^{-3}$

NA48 [PLB 604(2004), HEP2005 289]

K^0_{e3} : No evidence for quadratic term

$K^0_{\mu 3}$ linear fit: $\lambda_0 = (12.0 \pm 1.7) 10^{-3}$

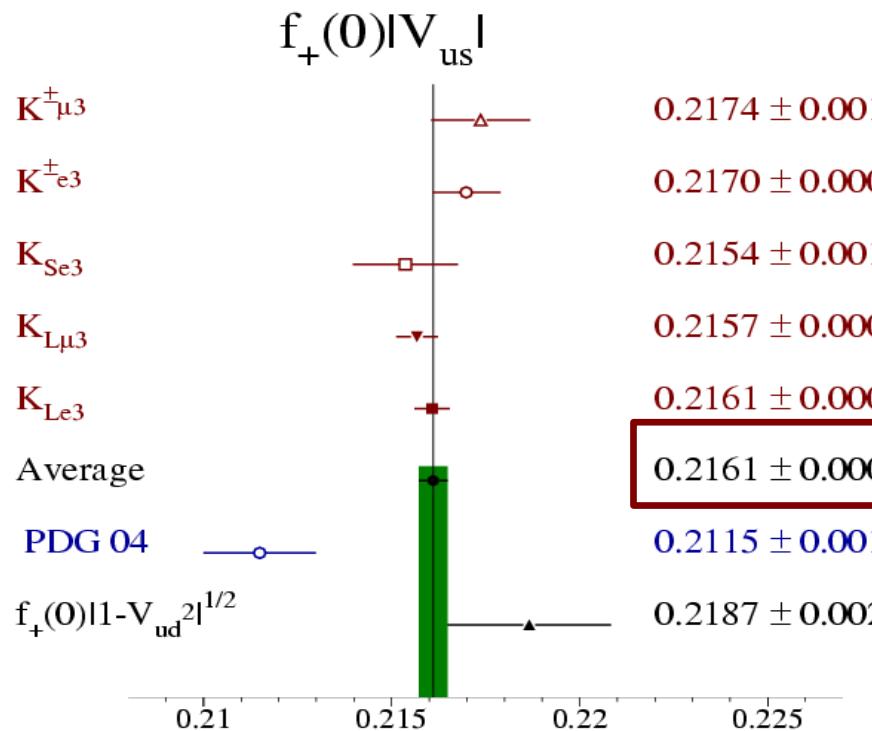


Experiment

V_{us} determination from Kl3

	$K_L e3$	$K_L \mu 3$	$K_S e3$	$K^\pm e3$	$K^\pm \mu 3$
BR	0.4046(8)	0.2697(7)	$7.046(91) \times 10^{-4}$	0.05043(31)	0.03383(27)
τ	51.10(19) ns		89.58(6) ps		12.386(22) ns

$$\begin{aligned}\lambda'_+ &= 0.02496(79) \\ \lambda''_+ &= 0.0016(3) \\ \lambda_0 &= 0.01587(95)\end{aligned}$$



- $f_+(0)=0.961(8)$
Leutwyler and Roos Z.
[Phys. C25, 91, 1984]
- $V_{ud}=0.97377(27)$
Marciano and Sirlin
[Phys.Rev.Lett.96 032002,2006]

$$V_{us} = 0.2249 \pm 0.0019$$

From $R_{e\mu}$:

$$\frac{g_e}{g_\mu} = 1.0014 \pm 0.0037$$

Experiment

$V_{us} - V_{ud}$ plane

Combining the experimental value of $\Gamma(K \rightarrow \mu\nu(\gamma))/\Gamma(\pi \rightarrow \mu\nu(\gamma))$

with the ratio f_K/f_π obtained from lattice calculations we can extract $|V_{us}|/|V_{ud}|$

(Marciano hep-ph/0406324) $\Gamma(K \rightarrow \mu\nu(\gamma))/\Gamma(\pi \rightarrow \mu\nu(\gamma)) \propto |V_{us}|^2/|V_{ud}|^2 f_K^2/f_\pi^2$

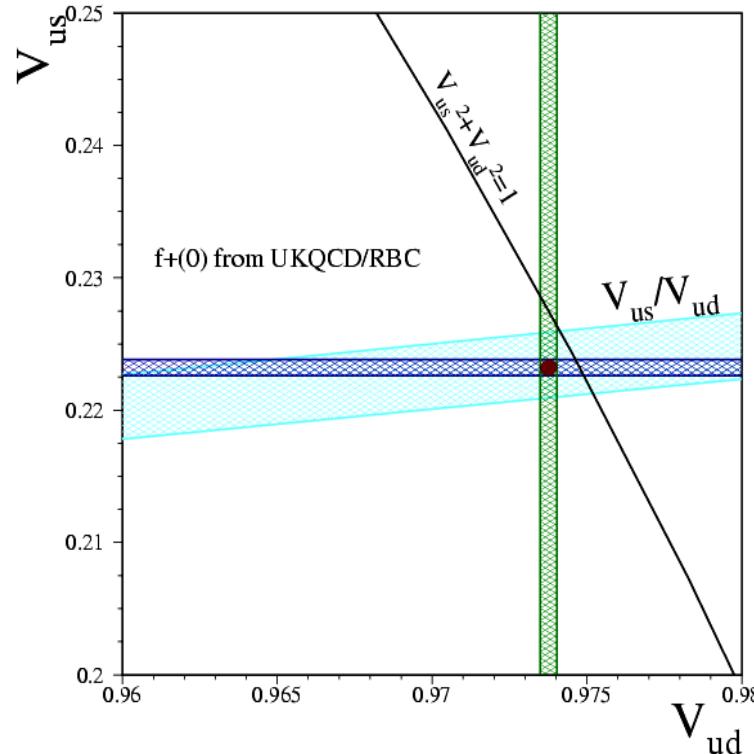
Using $f_K/f_\pi = 1.198(3)(^{+16}_{-5})$ from MILC

and KLOE BR($K^+ \rightarrow \mu^+\nu$)

we get $V_{us}/V_{ud} = 0.2294 \pm 0.0026$

$$\delta = |V_{us}|^2 + |V_{us}|^2 - 1 = -0.0020 \pm 0.0006$$

NEED TO BE CONFIRMED



A WG on Precise SM tests in K decays M.Antonelli LNF-INFN – FLAVIANET meeting Barcelona

Conclusions

- Chiral Perturbation Theory doing fine
- $\eta \rightarrow 3\pi$: All three have expressions: do we agree?
- $K_{\ell 3}$: useful even with very many parameters
- Other recent work: connection with lattice QCD:
Partially quenched ChPT
 - PQChPT at two loops: basic calculations done
 - Photons now also included (but NLO)
- For formulas:
<http://www.thep.lu.se/~bijnen/chpt.html>
- Many more calculations have been done