

Topological properties of citation and metabolic networks

Sven Bilke* and Carsten Peterson†

Complex Systems Division, Department of Theoretical Physics, University of Lund, Sölvegatan 14A, S-223 62 Lund, Sweden

(Received 19 March 2001; published 20 August 2001)

Topological properties of “scale-free” networks are investigated by determining their spectral dimensions d_S , which reflect a diffusion process in the corresponding graphs. Data bases for citation networks and metabolic networks together with simulation results from the growing network model [A.-L. Barabasi and R. Albert, *Science* **286**, 509 (1999)] are probed. For completeness and comparisons lattice, random and small-world models are also investigated. We find that d_S is around 3 for citation and metabolic networks, which is significantly different from the growing network model, for which d_S is approximately 7.5. This signals a substantial difference in network topology despite the observed similarities in vertex-order distributions. In addition, the diffusion analysis indicates that the citation networks are treelike in structure, whereas the metabolic networks contain many loops.

DOI: 10.1103/PhysRevE.64.036106

PACS number(s): 05.10.-a, 05.40.-a, 05.50.+q, 87.18.Sn

I. INTRODUCTION

There has recently been an upsurge of interest in the so-called scale-free networks, where the vertex order, or degree of connectivity per node k , follows a power-law distribution

$$P(k) \sim k^{-\gamma} \quad (1)$$

for $k > \langle k \rangle$. This is in contrast to the exponential suppression expected from randomly wired networks. For a number of real-world networks like social networks, power grids, citation networks, the world wide web, and metabolic networks [1], the scale-free behavior of Eq. (1) has been observed with exponents γ in the range 2–4 [2], at least within an appropriate range of orders [3]. The nature of all examples above is that they originate from a growing process. Hence, it appears natural to develop a model of growing network, which was indeed done in [2].

The degree of connectivity is a measure of local properties of the networks. Less attention has been paid to more global properties for the above examples except for general features like network diameters. Global measures could be important discriminants of network properties, in particular, when it comes to narrow on potential underlying models.

In this paper, we analyze topological properties of networks, both of synthetic and real-world nature, by extracting spectral dimensions d_S using a random walk procedure, from which the return-to-origin probability is estimated. The focus is on scale-free networks; citation networks, metabolic networks, and the growing network model [2] are investigated. For completeness and comparisons, we also extract the d_S from simple three-dimensional regular lattice networks, random networks, and small-world networks [4].

We find that the citation and metabolic networks are quite low dimensional with d_S around 3, whereas for the growing network model d_S is approximately 7.5. The latter turns out to be more in parity with what characterizes random and

small-world networks. As a consistency check, the dimension of regular lattice networks is also determined, which comes out as expected. One concludes that the spectral dimension offers a powerful and additional measure to γ in Eq. (1) when it comes to characterizing network topologies. Furthermore, the diffusion process underlying the extraction of d_S hints upon differences in the citation and metabolic network topologies; the former is treelike whereas the latter is rich in loop structures.

This paper is organized as follows: In Sec. II we describe the method for extracting the spectral dimension and Sec. III, IV, and V contain our investigations of the synthetic, citation, and metabolic networks, respectively. A summary can be found in Sec. VI.

II. SPECTRAL DIMENSION

Our method to probe the topological properties of an interaction network, which is of more global nature than the degree of connectivity, is based upon the diffusion of a test particle in the metric space defined by the graph. In a continuous space with a fixed smooth metric, the diffusion equation has the form

$$\frac{\partial}{\partial t} K_g(\epsilon, \epsilon_0, t) = \Delta_g K_g(\epsilon, \epsilon_0, t). \quad (2)$$

Here t is the diffusion time, Δ_g is the Laplace operator in the metric g , and $K_g(\epsilon, \epsilon_0, t)$ is the probability density to diffuse from ϵ_0 to ϵ in time t . For small t it is well known [5], that the average *return probability* has the following asymptotic expansion

$$K_g(0,0,t) \sim t^{-d_S/2} \quad (3)$$

with

$$K_g(\epsilon,0,0) = \delta(\epsilon) \quad (4)$$

and where the power $t^{-d_S/2}$ reflects the dimension of the network.

*Email address: sven@thep.lu.se

†Email address: carsten@thep.lu.se

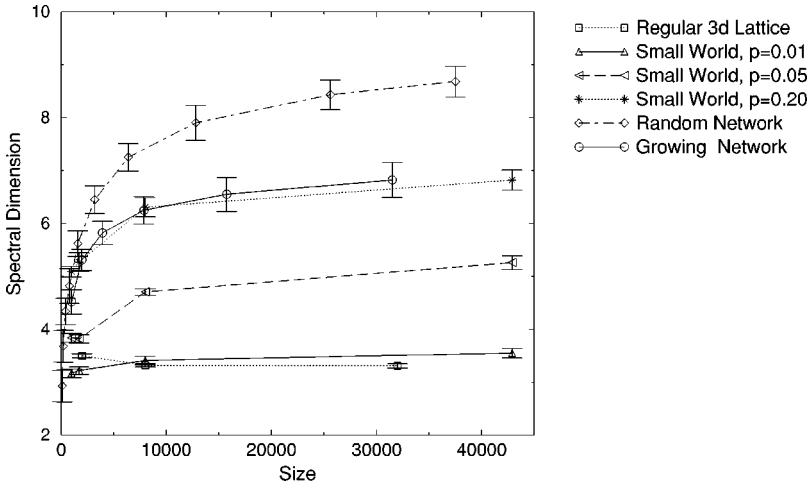


FIG. 1. Spectral dimensions d_S as functions of size for synthetic networks.

In the spirit of these equations we extract the *spectral dimension* of the geometry defined by the interaction networks considered in this work. To this end, we use the transition matrix

$$C_{ij} = \frac{J_{ij}}{k_i}, \quad (5)$$

with

$$J_{ij} = \begin{cases} 1 & \text{if } i \text{ and } j \text{ are neighbors} \\ 0 & \text{otherwise} \end{cases} \quad (6)$$

as the discrete version of the Laplace operator Δ_g . The number k_i , counts the number of links connected to vertex i —the vertex order. In a simulation of the diffusion process defined by the transition matrix, the probability to return, P , after t steps to the origin δ_i is then measured. More precisely, we choose a random subset $\{\delta_i\}, i=1 \dots N$ of vertices and extract

$$\langle P(t) \rangle = \frac{1}{N} \sum_i^N \langle \delta_i | C_{i,j}^T | \delta_i \rangle. \quad (7)$$

We then fit the resulting distribution to the asymptotic form in Eq. (3). For large t , Eq. (3) is dominated by the eigenvector of C_{ij} with eigenvalue $\lambda = 1$; the diffused particles reach an equilibrium distribution, and Eq. (3) does not hold. Also, for too small t the assumption of a smooth metric is not justified. Our extraction procedure is therefore somewhat more elaborate to account for these effects. A sliding window method is used, where the window is chosen such as to minimize the standard deviation per point used in the fit. In other words, we fit to the part of the distribution, which is closest to the functional form assumed in this procedure. Depending upon d_S , window sizes might vary from a few time points to the entire range.

III. SYNTHETIC NETWORKS

Next we extract d_S for a few different types of synthetic networks with sizes up to around 30 000 nodes in order to

compare with typical sizes of citation networks. In addition, for all networks the average geodesic distance (diameter) R is computed.

A. Lattice networks

We use a simple *square* lattices with dimension $d=3$. For these it is only possible to return to the origin after an even number of steps. Therefore $P(t)=0$ for odd t . In terms of the transfer matrix eigenvectors this means that the symmetry of the square lattice leads to an eigenvector with λ exactly -1 . The corresponding eigenstate does not decay for large t and yields destructive interference with $\lambda = 1$ eigenstate after an odd number of steps. When extracting the spectral dimension d_S from the data we therefore omit the odd time points. As can be seen from Fig. 1, d_S is slightly above the expected value 3, the dimension of the lattice. We have repeated this experiment for a regular lattice including diagonals, hence allowing for closed loops with an odd number of steps. Although local properties of the graph are considerably changed since the number of links emerging from each vertex is increased from $2d$ to $2d + d(d-1)$, we still extract the same value for d_S within errors. This is a further justification of omitting the data at odd time points for the square lattice and emphasizes that the spectral dimension is sensitive to global geometric aspects rather than to local details.

B. Small-world networks

We generate small-world networks [4] from the lattice network described above by rewiring the edges in the regular lattice to a randomly chosen vertex with probability $P = 0.01, 0.05, \text{ and } 0.2$ respectively. This generates three sets of models for each system size. Not surprisingly, these networks end up with spectral dimensions (see Fig. 1) in between the regular lattice and the homogeneous random networks described next.

C. Random networks

Homogeneous random networks are generated as directed networks, where each vertex has two inputs from other vertices chosen at random. In the analysis the graph is consid-

ered undirected by ignoring the orientation on the edges connecting the vertices. Furthermore, the analysis is focused on the largest cluster; the largest connected part of the resulting graph. The spectral dimension $d_S \approx 8.5$ (see Fig. 1) is relatively large. This agrees qualitatively with the observation that a characteristic length scale, e.g., average distance between vertices along links, grows slowly with the volume. It also signals a large dimensionality for this type of graph. It is interesting to note that again a distinct even-odd disparity emerges; the amplitude after an odd number of steps is almost zero. This indicates that the random network is dominated by a treelike structure, the number of loops (at least of odd length) is negligibly small compared to the amplitude obtained from backtracking the same path.

D. Growing network model

In this model [2] one at each time step t adds a node, which connects to m existing nodes (i) with a probability Π_i given by

$$\Pi_i = \frac{k_i^\beta}{\sum_j k_j^\beta}, \quad (8)$$

where k_i is the connectivity of node i and β is a parameter. In [2] the $\beta=1$ case was investigated both analytically in the $t \rightarrow \infty$ case and numerically for finite t . With the approximations involved one arrived at $\gamma=3$ in the analytic approach and $\gamma=2.9 \pm 0.1$ from simulations. For these networks, R is approximately 6 for large system sizes.

IV. CITATION NETWORK

In a citation network the links and nodes correspond to the citations and the publications (citing and cited), respectively. We use the SPIRES data base [6] for our citation network studies. This data base is limited to high energy physics publications but is not confined to articles that have been published in referred journals—citations to and from conference reports are also present. From the data we construct citation networks, which we treat as undirected, not fully connected graphs. By considering graphs generated by publications in a certain year or time span and the cited papers (regardless of publication year) we obtain a whole set of graphs. In these not fully connected graphs we focus on the largest connected cluster, the corresponding sizes are shown in Table I. When computing the connectivity distribution of the nodes, we confirm the power-law distribution [cf. Eq. (1)] already observed in citation networks [8] with $\gamma=2.7$.

In Fig. 2 we plot the probability to return to the origin after t time steps for a few typical networks. It is interesting to note that, within errors, the return probability for the 76 and 79 networks are in very good agreement indicating universal topological properties. This holds true also when comparing all of the other one-year networks with the exception of the 75 and the 87 networks (not shown), which cannot be fitted to Eq. (3). In these years the connectivity does not probably generate a homogeneous metric, because the geom-

TABLE I. Largest cluster sizes in the SPIRES data base for the years 1975–1989.

Year	Largest cluster	Year	Largest cluster
1975	20931	1983	32752
1976	22969	1984	34558
1977	23936	1985	37020
1978	26038	1986	39962
1979	27055	1987	44392
1980	28045	1988	47290
1981	29309	1989	45549
1982	31516	1975–1981	98104

etry is dominated by two or more large clusters, which are interconnected by only a few links.

For the networks composed by several years we again observe a remarkable universality when comparing different time periods. However, the distribution observed is considerably different from the one-year distributions. The reason presumably is a maturation of the citation networks leading to a tighter interconnection of the central cluster. In turn we observe larger dimension as can be seen from the steeper slope of the corresponding graph in Fig. 2. These issues for the citation networks have been subject to further investigations [9].

V. METABOLIC NETWORK

Another real-world interaction network is the metabolic network found in living cells. In these networks, substrates are treated as vertices, while chemical reactions connecting substrates and educts are treated as directed links. Recently it has been demonstrated that for this type of networks the

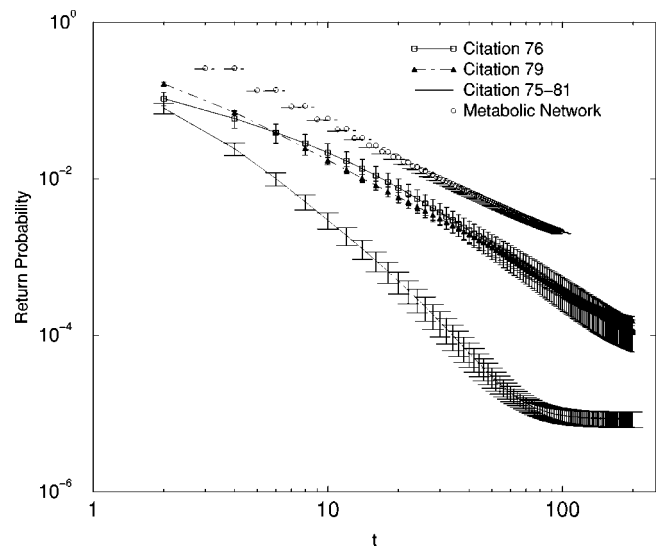


FIG. 2. The return probability (P) as a function of time steps t for citation networks with citing publications from a single years, 1976 and 1979, respectively, and for a six-year period (1976–1981). Same quantity for the metabolic network. To maintain readability, P for the metabolic network is multiplied with a factor 5.

TABLE II. Largest cluster sizes (N), average distances (diameters) (R), and spectral dimensions (d_S) for different network models and real-world networks. Also indicated are whether the return probabilities (P) are substantial for even or odd steps.

Network	N	R	d_S	$P_{\text{odd}}/P_{\text{even}}$
Regular lattice (d=3)	32000	14.3	3.1	0 (1 with diagonal)
Random network	32000	8.2	≈ 8.5	0
Small-world network ($p=0.01$)	32000	5.9	3.6	0
Small-world network ($p=0.05$)	32000	5.7	5.3	0
Small-world network ($p=0.2$)	32000	5.3	6.8	0
Network growing model	32000	5.9	≈ 7	0
Citation network	20000–200000	6.3–5.6	2.8–4.2	0
Metabolic network	3800	3.1	2.8	1

connectivity distribution obeys a power-law behavior [cf. Eq. (1)]; these networks are scale free with respect to the order distribution [1].

Using data from the EMP Project [7], we constructed a network including all reactions found in the database without taking into account species and cell locations. As above, we neglect the orientation of the resulting graph in the analysis. The spectral dimension observed for this network is $d_S \approx 2.8$. This is a surprisingly small dimension taking into account the observation [1] that the average distance between vertices on this graph does not grow with the graph size, implying a very large, possibly infinite dimension. Also, comparing d_S with the corresponding value for the growing model the scale-free network are analyzed above, the difference is remarkable. This indicates that, although the metabolic network and the scale-free network are similar on a local (vertex order) scale, the more global topological properties are very different. Another important difference arises from the observation that the return probability for odd path lengths do not vanish. This means that the metabolic networks consist of a very large number of closed loops. The graph has, in contrast to the random, growing model and citation networks, *not* a dominantly treelike structure. We interpret this as an indication of a built-in stability of metabolic networks with respect to small modifications.

VI. SUMMARY AND OUTLOOK

In Table II the results for d_S are summarized for the dif-

ferent networks. Also shown here are the average (geodesic) distances R . For the regular lattice, the spectral dimension reproduces quite well the dimension of the underlying lattice. Depending upon the probability to rewire a connection (p), the small-world networks takes on values between the regular lattice ($p=0$) and the random network ($p=1$), which has the largest spectral dimension of all networks probed here. One might have guessed that the growing network has an even larger dimension—the average distance for a network with the same size is smaller. However, in contrast to the geodesic distance, which only counts the shortest paths, the spectral dimension takes into account all possible paths of a given length. This means that only a few links, in extreme cases even a single link, can considerably change the geodesic distance, while the spectral dimension probes larger parts of the geometry and is therefore only slightly affected. For example, consider the behavior of the average distance for small-world networks with different probabilities P to rewire the underlying three-dimensional regular structure (Table II). Even for $P=0.01$ with the resulting geometry still essentially a regular three-dimensional lattice, the geodesic distance jumps to less than one-half of the value for the regular case, while the spectral dimension is changed only moderately.

The spectral dimensions for the real-world networks, the citation and metabolic networks, respectively, are strikingly similar (see Fig. 3) in view of the much larger values observed for the network growing model and the random network. We interpret this as a sign of universality in the ge-

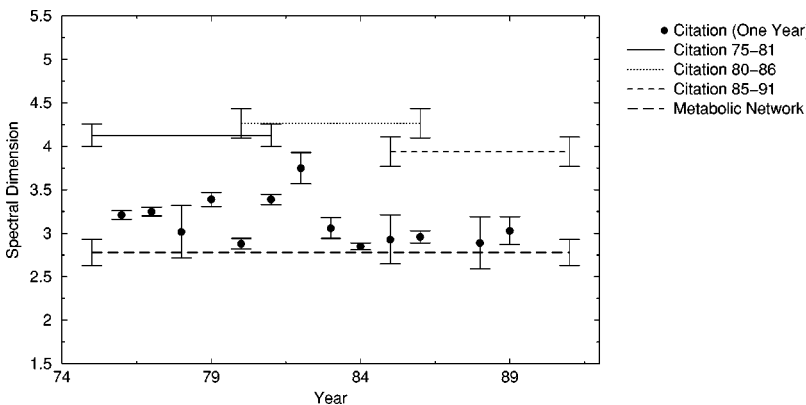


FIG. 3. The spectral dimension d_S for the citation networks for papers citing in one year and over six-year periods. For comparison, d_S for the metabolic network is also shown.

ometry of the real-world networks. The difference in d_S when comparing to that of the network growing model is remarkable. As mentioned in the Introduction, the latter reproduces quite nicely the power-law scale-free vertex-order distribution observed for the real-world networks [2]. However, the difference in d_S observed here indicates that the geometries of real-world networks are not fully described by the network growing model.

Given the similarity in d_S for the citation and metabolic networks, one might ask how similar they really are? An important difference can be observed by investigating the probability P_{odd} to return to the origin after an odd number of steps. For the citation network, as for all the synthetic networks, one obtains $P_{odd} \approx 0$. This is in contrast to the metabolic network for which $P_{odd} \approx P_{even}$. The case $P_{odd} \approx 0$ indicates a treelike geometrical structure with only few loops, where the return probability is dominated by the inverse of the forward paths. The regular lattice without diagonals seems to be a counter example. The return probability P_{odd} is exactly zero, but the geometry is certainly not tree-like and there exist many nonbacktracking loops. However, this lattice is a special case since its symmetry only allows for even step nonbacktracking loops. It is extremely unlikely that such an effect plays a role in any of the other lattices.

Our results can be summarized as follows.

(1) The spectral dimension d_S , which reflects a diffusion process on networks, offers quite some promise in categorizing networks beyond what emerges from studying vertex-order distributions.

(2) In particular, we find that d_S are very similar for citation and metabolic networks. This may indicate a universal behavior of real-world networks.

(3) With respect to d_S , the growing network model is significantly different from citation and metabolic networks.

(4) As a by-product when extracting d_S , we find that citation networks have treelike structures whereas the metabolic networks appear to contain many loops.

In addition, lattice, random, and small-world networks were probed using d_S with results exhibiting internal consistency of the method.

When finishing the write up of this work, a paper with somewhat similar scope and philosophy but with employing eigenvalue analysis techniques instead was released [10].

ACKNOWLEDGMENTS

We are very much indebted to Louise Addis in the SPIRES-HEP data base group at the Stanford Linear Accelerator Center library for providing us with the SPIRES data base in a form that substantially facilitated this project. Alex Selkov from the EMP-Project supported us in acquiring the data for the metabolic network. We have also benefited from discussions with Bo Söderberg regarding eigenvalue properties of network diffusion. This work was in part supported by the Swedish Foundation for Strategic Research and the Knut and Alice Wallenberg Foundation through the SWEGENE consortium.

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