Monte Carlo Event Generators

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1. (Monday) Introduction and Overview; Monte Carlo Techniques
   2. (Monday) Matrix Elements; Parton Showers I
   3. (yesterday) Parton Showers II; Matching Issues
   4. (yesterday) Multiple Interactions and Beam Remnants
   5. (today) Hadronization and Decays; Summary and Outlook
Event Physics Overview

Repetition: from the “simple” to the “complex”,
or from “calculable” at large virtualities to “modelled” at small

Matrix elements (ME):

1) Hard subprocess:
$|\mathcal{M}|^2$, Breit-Wigners, parton densities.

2) Resonance decays:
includes correlations.

Parton Showers (PS):

3) Final-state parton showers.
- $q \rightarrow qg$
- $g \rightarrow gg$
- $q \rightarrow q\bar{q}$
- $q \rightarrow q\gamma$

4) Initial-state parton showers.

$q$

$g$
5) Multiple parton–parton interactions.

6) Beam remnants, with colour connections.

5) + 6) = Underlying Event

7) Hadronization

8) Ordinary decays: hadronic, τ, charm, ...
Hadronization/Fragmentation models

Perturbative → nonperturbative → not calculable from first principles!

Model building = ideology + “cookbook”

Common approaches:

1) **String** Fragmentation
   (most ideological)

2) **Cluster** Fragmentation
   (simplest?)

3) **Independent** Fragmentation
   (most cookbook)

4) Local Parton–Hadron Duality
   (limited applicability)

Best studied in
\[ e^+e^- \rightarrow \gamma^*/Z^0 \rightarrow q\bar{q} \]
The Lund String Model

In QED, field lines go all the way to infinity

since photons cannot interact with each other.

Potential is simply additive:

$$V(x) \propto \sum_i \frac{1}{|x - x_i|}$$
In QCD, for large charge separation, field lines seem to be compressed to tubelike region(s) ⇒ **string(s)**

by self-interactions among soft gluons in the “vacuum”.
(Non-trivial ground state with quark and gluon “condensates”.
Analogy: vortex lines in type II superconductor)

Gives linear confinement with string tension:
\[ F(r) \approx \text{const} = \kappa \approx 1 \text{ GeV/fm} \quad \Leftrightarrow \quad V(r) \approx \kappa r \]

Separation of transverse and longitudinal degrees of freedom
⇒ simple description as 1+1-dimensional object – string –
with Lorentz invariant formalism
Linear confinement confirmed e.g. by quenched lattice QCD

\[ V(r) \approx -\frac{4 \alpha_s}{3} \frac{1}{r} + \kappa r \approx -\frac{0.13}{r} + r \]

(for \( \alpha_s \approx 0.5 \), \( r \) in fm and \( V \) in GeV)

\( V(0.4 \text{ fm}) \approx 0 \): Coulomb important for internal structure of hadrons, not for particle production (?)
Real world (??, or at least unquenched lattice QCD) \implies \text{nonperturbative string breakings } gg \ldots \rightarrow q\bar{q}

\text{quenched QCD}

\text{full QCD}

\text{Coulomb part}

\text{simplified colour representation:}
Repeat for large system ⇒ *Lund model* which neglects Coulomb part:

\[
\frac{dE}{dz} = \frac{dp_z}{dz} = \frac{dE}{dt} = \frac{dp_z}{dt} = \kappa
\]

Motion of quarks and antiquarks in a $q\bar{q}$ system:

This gives a simple but powerful picture of hadron production (with extensions to massive quarks, baryons, ...).
How does the string break?

\[ m_{\perp q'} = 0 \]

\[ d = \frac{m_{\perp q}}{\kappa} \]

\[ m_{\perp q'} > 0 \]

String breaking modelled by tunneling:

\[ \mathcal{P} \propto \exp \left( -\frac{\pi m_{\perp q}^2}{\kappa} \right) = \exp \left( -\frac{\pi p_{\perp q}^2}{\kappa} \right) \exp \left( -\frac{\pi m_{\perp q}^2}{\kappa} \right) \]

1) common Gaussian \( p_{\perp} \) spectrum

2) suppression of heavy quarks \( u\bar{u} : d\bar{d} : s\bar{s} : c\bar{c} \approx 1 : 1 : 0.3 : 10^{-11} \)

3) diquark \( \sim \) antiquark \( \Rightarrow \) simple model for baryon production

Hadron composition also depends on spin probabilities, hadronic wave functions, phase space, more complicated baryon production, …

\( \Rightarrow \) “moderate” predictivity (many parameters!)
Fragmentation starts in the middle and spreads outwards:

but breakup vertices causally disconnected
⇒ can proceed in arbitrary order
⇒ left–right symmetry

\[ \mathcal{P}(1, 2) = \mathcal{P}(1) \times \mathcal{P}(1 \rightarrow 2) \]
\[ = \mathcal{P}(2) \times \mathcal{P}(2 \rightarrow 1) \]

⇒ Lund symmetric fragmentation function
\[ f(z) \propto (1 - z)^a \exp(-b m^2_\perp / z) / z \]
The iterative ansatz

\[
q_0, p_{\perp 0}, p_+ \rightarrow q_0 q_1, p_{\perp 0} - p_{\perp 1}, z_1 p_+
\]

and so on until joining in the middle of the event

\[
q_1 q_1, p_{\perp 1} - p_{\perp 2}, z_2 (1 - z_1) p_+
\]

\[
q_2 q_2, p_{\perp 2} - p_{\perp 3}, z_3 (1 - z_2)(1 - z_1) p_+
\]

Scaling in lightcone \( p_\pm = E \pm p_z \) (for \( q\bar{q} \) system along \( z \) axis) implies flat central rapidity plateau + some endpoint effects:

\[
\frac{dn}{dy}
\]

\[
\langle n_{ch} \rangle \approx c_0 + c_1 \ln E_{cm}, \sim \text{Poissonian multiplicity distribution}
\]
The most characteristic feature of the Lund model

Gluon = kink on string, carrying energy and momentum

Force ratio gluon/ quark = 2, cf. QCD \( N_C/C_F = 9/4 \), \( \rightarrow 2 \) for \( N_C \rightarrow \infty \)

No new parameters introduced for gluon jets!, so:

- Few parameters to describe energy-momentum structure!
- Many parameters to describe flavour composition!
Independent fragmentation

Based on a similar iterative ansatz as string, but

\[ q = q + g + \bar{q} \]

minor corrections in middle

String effect (JADE, 1980)
\[ \approx \text{coherence in nonperturbative context} \]

Further numerous and detailed tests at LEP favour string picture . . .
. . . but much is still uncertain when moving to hadron colliders.
Lund news: fragmentation of junction topology

Encountered in $R$-parity violating SUSY decays $\tilde{\chi}_1^0 \to uds$, or when 2 valence quarks kicked out of proton beam

More complicated (but $\approx$ solved) with gluon emission and massive quarks
“Preconfinement”: colour flow is local in coherent shower evolution

1) Introduce forced $g \rightarrow q\bar{q}$ branchings
2) Form colour singlet clusters
3) Clusters decay isotropically to 2 hadrons according to phase space weight $\sim (2s_1 + 1)(2s_2 + 1)(2p^*/m)$

simple and clean, but …
1) Tail to very large-mass clusters (e.g. if no emission in shower); if large-mass cluster $\rightarrow$ 2 hadrons then incorrect hadron momentum spectrum, crazy four-jet events $\implies$ split big cluster into 2 smaller along "string" direction; daughter-mass spectrum $\Rightarrow$ iterate if required; $\sim 15\%$ of primary clusters are split, but give $\sim 50\%$ of final hadrons

2) Isotropic baryon decay inside cluster $\implies$ splittings $g \rightarrow q\bar{q} + q\bar{q}$

3) Too soft charm/bottom spectra $\implies$ anisotropic leading-cluster decay

4) Charge correlations still problematic $\implies$ all clusters anisotropic (?)

5) Sensitivity to particle content $\implies$ only include complete multiplets
String vs. Cluster

<table>
<thead>
<tr>
<th>Program Model</th>
<th>PYTHIA String</th>
<th>HERWIG Cluster</th>
</tr>
</thead>
<tbody>
<tr>
<td>Energy–Momentum Picture</td>
<td>Powerful, Predictive, Few</td>
<td>Simple, Unpredictive, Many</td>
</tr>
<tr>
<td>Parameters</td>
<td>Messy, Unpredictive, Many</td>
<td>Simple, In-between, Few</td>
</tr>
</tbody>
</table>

“There ain’t no such thing as a parameter-free *good* description”
Local Parton–Hadron Duality

Analytic approach:
Run shower down to to $Q \approx \Lambda_{QCD}$
(or $m_{\text{hadron}}$, if larger)

“Hard Line”: each parton $\equiv$ one hadron

“Soft Line”: local hadron density
$\propto$ parton density
describes momentum spectra $dn/dx_p$
and semi-inclusive particle flow,
but fails for identified particles

+ “renormalons” (power corrections)

\[
\langle 1 - T \rangle = a \alpha_s(E_{cm}) + b \alpha_s^2(E_{cm}) 
+ c/E_{cm}
\]

Not Monte Carlo, not for arbitrary quantities
Decays

Unspectacular/ungrateful but necessary:
this is where most of the final-state particles are produced!
Involves hundreds of particle kinds and thousands of decay modes.

e.g.

\[ B^*0 \rightarrow B^0 \gamma \]
\[ B^0 \rightarrow \bar{B}^0 \]
\[ D^{*+} \rightarrow \pi^+ e^- \]
\[ \rho^+ \rightarrow \pi^+ \pi^0 \]
\[ \pi^0 \rightarrow e^+ e^- \gamma \]

- \( B^*0 \rightarrow B^0 \gamma \): electromagnetic decay
- \( B^0 \rightarrow \bar{B}^0 \) mixing (weak)
- \( \bar{B}^0 \rightarrow D^{*+} \nu_e e^- \): weak decay, displaced vertex, \( |M|^2 \propto (p_{\bar{B}} p_{\nu})(p_e p_{D^*}) \)
- \( D^{*+} \rightarrow D^0 \pi^+ \): strong decay
- \( D^0 \rightarrow \rho^+ K^- \): weak decay, displaced vertex, \( \rho \) mass smeared
- \( \rho^+ \rightarrow \pi^+ \pi^0 \): \( \rho \) polarized, \( |M|^2 \propto \cos^2 \theta \) in \( \rho \) rest frame
- \( \pi^0 \rightarrow e^+ e^- \gamma \): Dalitz decay, \( m(e^+ e^-) \) peaked

Dedicated programs, with special attention to polarization effects:

- EVTGEN: \( B \) decays
- TAUOLA: \( \tau \) decays
Question: are jets the same in all processes?

Answer 1: no, at LEP mainly quarks jets, often $b/c$, at LHC mainly gluons, if quarks then mainly $u/d$.

Answer 2: no, perturbative evolution gives calculable differences.
Answer 3: (string) hadronization mechanism assumed universal, but is not quite.

\[ E \frac{d^3\sigma}{d^3p} : \text{Dependence on proton } P_T \]

so discrepancies

\[ \frac{\mathcal{P}_{qq}}{\mathcal{P}_q} = 0.1 \text{ at LEP}, \quad \frac{\mathcal{P}_s}{\mathcal{P}_u} = 0.3 \text{ at LEP} \]

\[ = 0.05 \text{ at HERA} \]

Reasons? HERA dominated by “beam jets”, so

- Less perturbative evolution ⇒ strings less “wrinkled”?
- Many overlapping strings ⇒ collective phenomena?
Other program tasks/elements

- Diffractive physics (≈ rapidity-gap physics)
  \[ \sigma_{\text{el}} \approx 25 \text{ mb} \quad \text{pp} \rightarrow \text{pp} \]
  LHC: \[ \sigma_{\text{diff}} \approx 25 \text{ mb} \quad \text{pp} \rightarrow \text{p}X, \text{pp} \rightarrow X_1X_2, \text{etc} \]
  \[ \sigma_{\text{inel,nondiff}} \approx 50 \text{ mb} \quad \text{pp} \rightarrow X \text{ (without obvious subdivision of } X) \]

\[ \frac{dn}{dy} \]

- Colour reconnection: how well can we trust “perturbatively” calculable colour flow in soft region?

- Bose-Einstein: must we use amplitudes to describe production of identical particles? (≈ 50 \( \pi^+ \), 50 \( \pi^- \), 70 \( \pi^0 \) per event)

- Event measures; jet clustering routines; other utilities

… and more
Event Generator Practicalities
Event generation structure

1) Initialization step
   - select process(es) to study
   - modify physics parameters: $m_t, m_h, \ldots$
     - set kinematics constraints
   - modify generator performance
     - initialize generator
     - book histograms

2) Generation loop
   - generate one event at a time
   - analyze it (or store for later use)
     - add results to histograms
     - print a few events

3) Finishing step
   - print deduced cross-sections
   - print/save histograms etc.
How to run event generators

Often forced to use what is allowed by constricted collaboration framework, but for maximal power and minimal bugs run raw generator:

- **HERWIG, ISAJET**: supplied but modifiable main program, calling user-written routines

  - **HWIGPR**: main program
    - Supplied, but needs modifying to initialize parameters, steer event generation, etc
  
  - **HERWIG**: subroutine library
    - Shouldn’t need modifying!

  - **HWABEG**: analysis initialization
  - **HWANAL**: event analysis
  - **HWAEND**: terminate analysis

- **PYTHIA**: generator is subroutine package, user writes main program
C... Arithmetic in double precision; integer functions; PYDATA.
IMPLICIT DOUBLE PRECISION(A-H, O-Z)
INTEGER PYK, PYCHGE, PYCOMP
EXTERNAL PYDATA

C... The event record and other common blocks.
COMMON PYJETS/N, NPAD, K(4000,5), P(4000,5), V(4000,5)
COMMON PYDAT2/KCHG(500,4), PMAS(500,4), PARF(2000), VCKM(4,4)
COMMON PYSUBS/MSEL, MSELPD, MSUB(500), KFIN(2,-40:40), CKIN(200)
COMMON PYPARS/MSTP(200), PARP(200), MSTI(200), PARI(200)

C... Physics scenario.
MSEL = 0 ! Mix subprocesses freely
MSUB(102) = 1 ! g + g \rightarrow h_0
MSUB(123) = 1 ! f + f' \rightarrow f + f' + h_0
MSUB(124) = 1 ! f + f' \rightarrow f'' + f'' + h_0
PMAS(25,1) = 300D0 ! Nominal Higgs mass.

C... Run parameters.
NEV = 1000 ! Number of events
ECM = 14000D0 ! CM energy of run
CKIN(1) = 200D0 ! Minimum Higgs mass.
CKIN(2) = 400D0 ! Maximum Higgs mass.

C... Initialize and book histogram(s).
CALL PYINIT('CMS', 'p', 'p', ECM)
CALL PYBOOK(1, 'Higgs mass distribution', 80, 200D0, 400D0)

C... Generate events and look at first few.
DO 200 IEV = 1, NEV
   CALL PYEVNT
   IF (IEV.LE.1) CALL PYLIST(1)
C... Find Higgs and fill its mass. End event loop.
DO 150 I = 7, 9
   IF (K(I,2).EQ.25) CALL PYFILL(1, P(I,5), 1D0)
150 CONTINUE
200 CONTINUE

C... Final output.
CALL PYSTAT(1) ! Print cross section table
CALL PYHIST ! Print histogram(s)
END
On To C++

Currently HERWIG and PYTHIA are successfully being used, also in new LHC environments, using C++ wrappers

Q: Why rewrite?
A1: Need to clean up!
A2: Fortran 77 is limiting

Q: Why C++?
A1: All the reasons for ROOT, Geant4, . . .
   (“a better language”, industrial standard, . . .)
A2: Young experimentalists will expect C++
   (educational and professional continuity)
A3: Only game in town! Fortran 90

So far mixed experience:
- Conversion effort: everything takes longer and costs more
  (as for LHC machine, detectors and software)
- The physics hurdle is as steep as the C++ learning curve
C++ Players

PYTHIA7 project $\implies$ **ThePEG**
Toolkit for High Energy Physics Event Generation
(L. Lönnblad; S. Gieseke, A. Ribon, P. Richardson)

HERWIG++: complete reimplementation
(S. Gieseke, D. Grellscheid, A. Ribon, P. Richardson,
M.H. Seymour, P. Stephens, B.R. Webber; M. Bähr, M. Gigg,
K. Hamilton, S. Latunde-Dada, S. Plätzer, A. Sherstnev)

ARIADNE/LDC: to do ISR/FSR showers, multiple interactions
(L. Lönnblad; N. Lavesson)

SHERPA: partly wrappers to PYTHIA Fortran; has CKKW
(F. Krauss; T. Fischer, T. Gleisberg, S. Hoeche, T. Laubrich,
R. Matyszkwiewicz, A. Schaelicke, C. Semmling, F. Siegert,
S. Schumann, J. Winter)

PYTHIA8: restart to write complete event generator
(T. Sjöstrand, (S. Mrenna, P. Skands))
What is ThePEG?

Toolkit for High Energy Physics Event Generation

CLHEP utilities

ThePEG basic structure

HERWIG++ physics modules

PYTHIA7 physics modules

Ariadne/LDC physics modules

? PYTHIA8 ?

not SHERPA
Running ThePEG

• ThePEG defines a set of abstract Handler classes for hard partonic sub-processes, parton densities, QCD cascades, hadronization, ...

• These handler classes interacts with the underlying structure using a special Event Record and a pre-defined set of virtual functions.

• The procedure to implement e.g. a new hadronization model, is to write a new (C++) class inheriting from the abstract HadronizationHandler base class, implementing the relevant virtual functions.

• The end-user will use a setup program to be able to pick objects corresponding to different physics models to build up an EventGenerator which then can be run interactively or off-line, or as a special slave program e.g. for Geant4.

• The setup program is used to choose between a multitude of pre-defined generators, to modify parameters and options of the selected models and, optionally, to specify the analysis to be done on the generated events.

• The Repository is the central part of the setup phase. It handles a structured list of all available objects and allows the user to manipulate them.
The new generator Herwig++

A completely new event generator in C++

- Aiming at full multi-purpose generator for LHC and future colliders.
- Preserving main features of HERWIG such as
  - angular ordered parton shower
  - cluster hadronization
- New features and improvements
  - covariant shower formulation
  - improved parton shower evolution for heavy quarks
  - consistent radiation from unstable particles (multiscale evolution)
Hard interactions

- Basic ME’s included in ThePEG, such as:

  \[ e^+ e^- \rightarrow q\bar{q}, \text{ partonic } 2 \rightarrow 2, \]

  we use them.

- Soft and hard matrix element corrections implemented for \( e^+ e^- \rightarrow q\bar{q}g \).

- AMEGIC++ will provide arbitrary ME’s for multiparton final states via AMEGICInterface.

- LesHouchesFileReader enables to read in and process any hard event generated by parton level event generators (MadGraph/MadEvent, AlpGen, CompHEP,...).

- CKKW ME+PS foreseen.

- Other authors can easily include their own matrix elements (→ safety of OO code)

**New/Future:** HELAS like structures are already implemented for decays and spin correlations allows us to code simple processes efficiently.
. . . and New Decays!

- Better decayers are being developed for almost all decay modes.
- \( \rightarrow B \) decays.
- Spin correlations will be included.
- Major effort ongoing
  - a universal database is being set up.
  - contains 448 particles and 2607 decay modes at present.
  - possibility to generate configuration files for different generators (they need to write their own code however . . ).
- Particle data book as guideline.

\( \rightarrow \text{look at examples. . . } \)
Herwig++ Particle Properties Database

This is the development version of the Herwig++ particle properties database. This is intended to replace the storage of particle properties as a text file to improve maintainability and accessibility.

This version is for the Herwig++ software only and much of the information is preliminary.

The database currently contains 448 particles and 2097 decay modes.

The information is stored in a number of forms:
- The particles numerically listed according to the PDG code
- The particles listed according to the multiplicity taken from the PDG
- The decays
- The Widths
- The Masses
- The Charges
- Generate .tea files for event generators

The contents of the database can be altered by following the links in the particle table or particle descriptions by selecting an option below:

- Add or modify a particle
- Add a decay mode for particle with ID
- Add a meson multiplet
- Add an overlap
- Add a mass generator
- Add a mass generator
- Add a Coupling
- Set the multiplets
- Set the decay modes for a particle to the charge conjugates of the anti-particle

Simple checks on the contents of the database:

- Check charge conservation in decays
- Check lepton number conservation in decays
- Check the spin is consistent with the PDG code

Last modified: Monday 31 17:56:08 GMT 2006

file://C://xy/gw/home/re/snow/TeXSN/Herwig++/index.html 15.06.2008
file://C://xy/gw/home/re/snow/TeXSN/Database/EnvDatabase.html 15.06.2008
What's next?

Near Future... 

★ Initial state shower:
  • Complete implementation and tests.
★ Refine $e^+e^-$:
  • Full CKKW ME+PS matching.
  • Precision tune to LEP data should be possible.
★ with IS and FS showers running:
  • we can start to test Drell–Yan and jets in pp collisions.
  • cross check with Tevatron data and finally make predictions for the LHC.
★ Underlying Event.
★ Hadronic Decays: **NEW!** many new decayers, $\tau$–decays, Spin correlations (P Richardson).
★ New Ideas: soft gluons, improved shower algorithm, NLO, . . .

Schedule?

• Ready for LHC!
Status of SHERPA

Scope:
- Full simulation of high energetic particle reactions at existing and future collider experiments, incl. $e^+e^-$, $\gamma\gamma$, $e\gamma$, $ep$, $p\bar{p}$ and $pp$ collisions

Method:
- Account for multi-jet production through tree level matrix elements
- Combine them with the parton showers and hadronization according to the CKKW prescription

First $\alpha$-version was released during MC4LHC workshop 2003, current version is SHERPA$\alpha$-1.0.7

Sources:
- downloads, manual, bug reports etc. under http://www.sherpa-mc.de
Conclusions/Outlook

SHERPA including the ME’s of AMEGIC++ and the CKKW prescription to combine them with the PS is a powerful tool to attempt the description of present-day Tevatron data and to study the extrapolation to LHC energies.

New features of version 1.0.7:
- Revised SUSY sector, including the SLHA interface
- $\tau$-lepton and first hadron decays
- Improved decay chain treatment

Things to do:
- Alternative underlying event model
- Include and tune the cluster fragmentation model
- Supplement the hadron decays, special emphasis on $b$-hadrons
- ...
PYTHIA8: A fresh start

Problem: PYTHIA7 stalled, no other manpower
Solution?: take a sabbatical and work “full-time”!
(⇒ baseline model, S. Mrenna & P. Skands join later ?)

Tentative schedule:

<table>
<thead>
<tr>
<th>time</th>
<th>date</th>
<th>processes</th>
<th>final states</th>
</tr>
</thead>
<tbody>
<tr>
<td>0</td>
<td>1 Sept. 2004</td>
<td>—</td>
<td>—</td>
</tr>
<tr>
<td>1</td>
<td>1 Sept. 2005</td>
<td>LHA-style input</td>
<td>incomplete draft</td>
</tr>
<tr>
<td>2</td>
<td>1 Sept. 2006</td>
<td>a few processes</td>
<td>complete, buggy(?)</td>
</tr>
<tr>
<td>3</td>
<td>1 Sept. 2007</td>
<td>more processes</td>
<td>stable, debugged</td>
</tr>
</tbody>
</table>

…but don’t forget Murphy’s law

Objectives:

- clean up, keep the most recent models
- Les Houches Accord style input central
- independent of ThePEG (or anything else), but
- interface to ThePEG later written by L. Lönnblad (?)
Current PYTHIA8 structure

The User (∼ Main Program)

Pythia

Event process

Event event

ProcessLevel

LHAinit
LHAevnt
(PYTHIA 6.3)
(...??)

PartonLevel

TimeShower
SpaceShower
MultipleInteractions
BeamRemnants

HadronLevel

StringFragmentation
MiniStringFrag...
ParticleDecays
(...??)

BeamParticle

Vec4, Random, Settings, ParticleData, StandardModel, ...
## Current PYTHIA8 status

<table>
<thead>
<tr>
<th>Existing classes</th>
<th>Missing classes/topics</th>
</tr>
</thead>
<tbody>
<tr>
<td><strong>Process Level</strong></td>
<td><strong>Cross section administration</strong></td>
</tr>
<tr>
<td>LHAinit **</td>
<td><strong>Phase space selection</strong></td>
</tr>
<tr>
<td>LHAevnt **</td>
<td><strong>Process matrix elements</strong></td>
</tr>
<tr>
<td>(PYTHIA 6.3) ***</td>
<td><strong>Parton density libraries</strong></td>
</tr>
<tr>
<td><strong>Parton Level</strong></td>
<td><strong>Resonance decays</strong></td>
</tr>
<tr>
<td>TimeShower **</td>
<td><strong>ThePEG input (?)</strong></td>
</tr>
<tr>
<td>SpaceShower **</td>
<td><strong>MI/ISR/FSR interleaving</strong></td>
</tr>
<tr>
<td>MultipleInteractions **</td>
<td><strong>colour flow models</strong></td>
</tr>
<tr>
<td>BeamRemnants **</td>
<td><strong>ME/PS matching</strong></td>
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<tr>
<td><strong>Hadron Level</strong></td>
<td><strong>Popcorn baryons</strong></td>
</tr>
<tr>
<td>StringFragmentation **</td>
<td>updated decay tables</td>
</tr>
<tr>
<td>MiniStringFrag... **</td>
<td>Bose-Einstein</td>
</tr>
<tr>
<td>ParticleDecays **</td>
<td>event analysis routines</td>
</tr>
<tr>
<td><strong>—</strong></td>
<td><strong>...and much, much more</strong></td>
</tr>
<tr>
<td>Event **</td>
<td></td>
</tr>
<tr>
<td>BeamParticle **</td>
<td></td>
</tr>
<tr>
<td>Vec4, Random ***</td>
<td></td>
</tr>
<tr>
<td>Settings **</td>
<td></td>
</tr>
<tr>
<td>ParticleData **</td>
<td></td>
</tr>
</tbody>
</table>
Outlook

Generators in state of continuous development:

☆ better & more user-friendly general-purpose matrix element calculators + integrators ☆
☆ new libraries of physics processes, also to NLO ☆
   ☆ more precise parton showers ☆
   ☆ better matching matrix elements ⇔ showers ☆
☆ improved models for underlying events / minimum bias ☆
   ☆ upgrades of hadronization and decays ☆
   ☆ moving to C++ ☆
⇒ always better, but never enough

But what are the alternatives, when event structures are complicated and analytical methods inadequate?
Final Words of Warning

[... ] The Monte Carlo simulation has become the major means of visualization of not only detector performance but also of physics phenomena. So far so good. But it often happens that the physics simulations provided by the Monte Carlo generators carry the authority of data itself. They look like data and feel like data, and if one is not careful they are accepted as if they were data.

[... ] I am prepared to believe that the computer-literate generation (of which I am a little too old to be a member) is in principle no less competent and in fact benefits relative to us in the older generation by having these marvelous tools. They do allow one to look at, indeed visualize, the problems in new ways. But I also fear a kind of “terminal illness”, perhaps traceable to the influence of television at an early age. There the way one learns is simply to passively stare into a screen and wait for the truth to be delivered. A number of physicists nowadays seem to do just this.

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