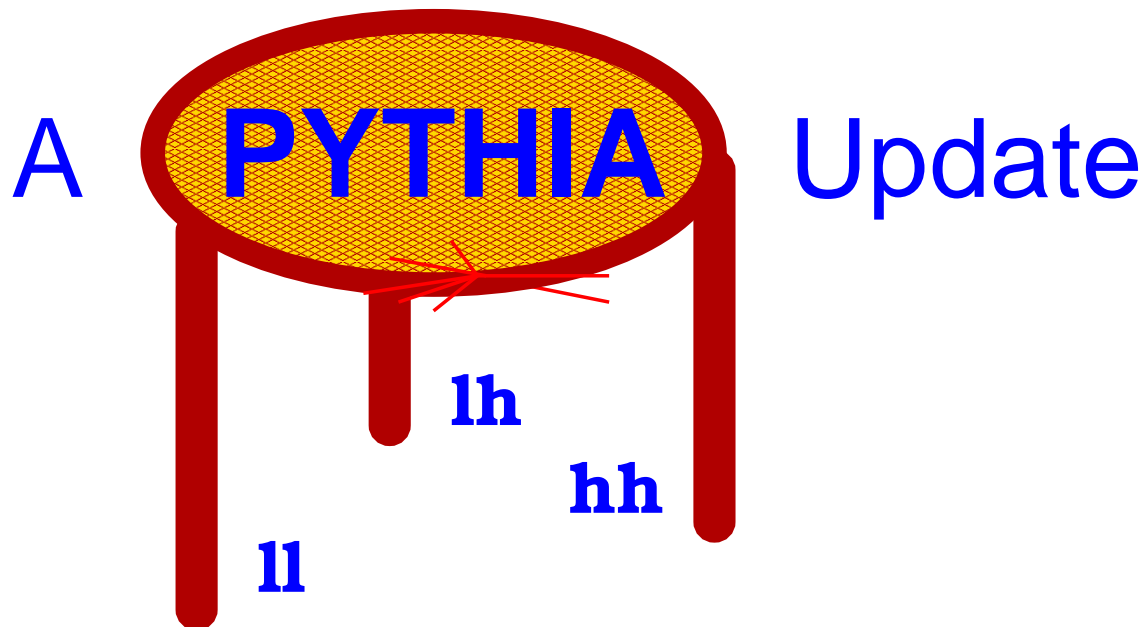


ECFA/DESY Linear Collider Workshop  
Oxford 20–23 March 1999



Torbjörn Sjöstrand  
Lund University

Introduction

Virtual Photon Processes

The photon flux • Photon processes

Parton distributions • Other model aspects

Some results • Program particulars

Photon ISR in  $Z^0$  production

Doubly-charged Higgses

Other news

On to C++

# Introduction



General-purpose generator:

- Hard subprocess “library”
- Convolution with parton distributions (cross sections, kinematics)
- Resonance decays (process-dependent)
- Initial-state QCD (& QED) showers
- Final-state QCD & QED showers
- Multiple parton–parton interactions
- Beam remnants
- Hadronization (string fragmentation)
- Decay chains
- Analysis & utility routines

Currently PYTHIA 6.125 of 21 February 1999

~ 46.800 lines Fortran 77

Code, manuals, sample main programs:

<http://www.thep.lu.se/~torbjorn/Pythia.html>

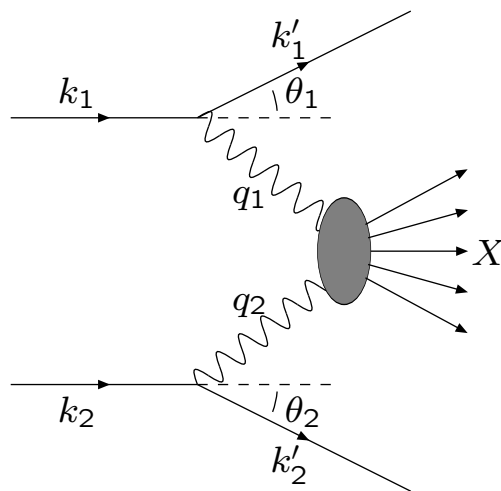
# PYTHIA 6.1 main news

- JETSET routines renamed:  
LUxxxx → PYxxxx + some more
- All real variables in DOUBLE PRECISION
- New SUSY processes and improved SUSY simulation; new PDG codes for sparticles
- New processes for Higgs, technicolour, ...
- Alternative Higgs mass shape
- Newer parton distributions (but ...)
- Several improved resonance decays
- Initial-state showers matched to (some) matrix elements
- QED radiation off an incoming muon
- New machinery to handle real and virtual photon fluxes and cross sections
- Energy-dependent  $p_{\perp\min}$  in multiple interactions
- Colour rearrangement options for  $W^+W^-$
- Expanded Bose-Einstein algorithm
- New baryon production scheme (optional)
- One-dimensional histograms (GBOOK)

# Virtual photon processes

(C. Friberg & TS, in preparation)

## The photon flux



$$W^2 = (q_1 + q_2)^2$$

$$Q_i^2 = -q_i^2$$

$$y_i = \frac{q_i k_j}{k_i k_j} \quad j=2(1) \text{ for } i=1(2)$$

$$x_i = \frac{q_i (k_1 + k_2)}{k_i (k_1 + k_2)}$$

$$y_i \approx x_i + \frac{Q_i^2}{s}$$

$$Q_i^2 \approx \frac{x_i^2}{1 - x_i} m_i^2 + (1 - x_i) s \sin^2(\theta_i/2)$$

$$d\sigma(ee \rightarrow eeX) = \sum_{\xi_1, \xi_2 = T, L} \int dy_1 dQ_1^2 dy_2 dQ_2^2$$

$$\times f_{\gamma/e}^{\xi_1}(y_1, Q_1^2) f_{\gamma/e}^{\xi_2}(y_2, Q_2^2) d\sigma(\gamma_{\xi_1}^* \gamma_{\xi_2}^* \rightarrow X)$$

with

$$f_{\gamma/e}^T(y, Q^2) = \frac{\alpha_{em}}{2\pi} \left( \frac{(1 + (1 - y)^2)}{y} \frac{1}{Q^2} - \frac{2m_e^2 y}{Q^4} \right)$$

$$f_{\gamma/e}^L(y, Q^2) = \frac{\alpha_{em}}{2\pi} \frac{2(1 - y)}{y} \frac{1}{Q^2}$$

# Photon processes

So far only jet production, i.e. not low- $p_{\perp}$ .

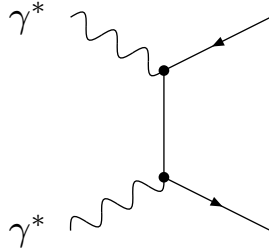
Three main process classes:

## 1. direct × direct

$$\gamma_T^* \gamma_T^* \rightarrow q\bar{q}$$

$$\gamma_T^* \gamma_L^* \rightarrow q\bar{q}$$

$$\gamma_L^* \gamma_L^* \rightarrow q\bar{q}$$



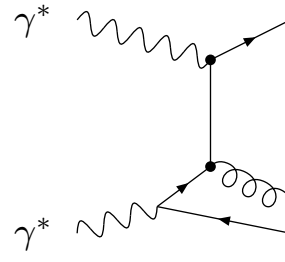
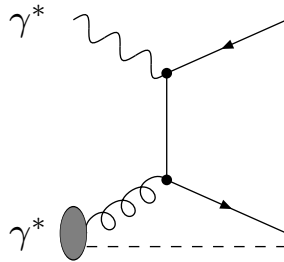
## 2. direct × resolved, resolved = VMD or anomalous

$$\gamma_T^* g \rightarrow q\bar{q}$$

$$\gamma_L^* g \rightarrow q\bar{q}$$

$$\gamma_T^* q \rightarrow qg$$

$$\gamma_L^* q \rightarrow qg$$



etc.

## 3. resolved × resolved

$$qq' \rightarrow qq'$$

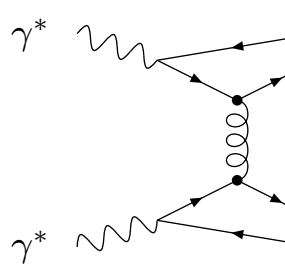
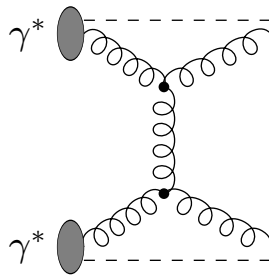
$$q\bar{q} \rightarrow q'\bar{q}'$$

$$q\bar{q} \rightarrow gg$$

$$qg \rightarrow qg$$

$$gg \rightarrow gg$$

$$gg \rightarrow q\bar{q}$$



etc.

$$d\sigma(\gamma^* \gamma^* \rightarrow X) = \left( \int d\hat{x}_1 f_i^{\gamma^*}(\hat{x}_1, \mu^2, Q_1^2) \right) \\ \times \left( \int d\hat{x}_2 f_j^{\gamma^*}(\hat{x}_2, \mu^2, Q_2^2) \right) \int d\hat{t} \frac{d\hat{\sigma}}{d\hat{t}}(\hat{s} = \hat{x}_1 \hat{x}_2 W^2)$$

## Parton distributions

SaS  $\gamma^*$  distributions extended from  $\gamma$  ones by inclusion of dipole damping factors,

$$f_a^{\gamma^*}(x, \mu^2, Q^2) = \sum_V \frac{4\pi\alpha_{em}}{f_V^2} \left( \frac{m_V^2}{m_V^2 + Q^2} \right)^2 f_a^{\gamma, V}(x, \mu^2; \tilde{Q}_0^2) + \frac{\alpha_{em}}{2\pi} \sum_q 2e_q^2 \int_{Q_0^2}^{\mu^2} \frac{dk^2}{k^2} \left( \frac{k^2}{k^2 + Q^2} \right)^2 f_a^{\gamma, q\bar{q}}(x, \mu^2; k^2),$$

reduction of evolution range:  $Q_0 < \tilde{Q}_0 < \mu$ ,  
and matching  $f_a^{\gamma^*}(x, \mu^2, Q^2) \rightarrow 0$  for  $\mu^2 \rightarrow Q^2$ .

$\mu^2$  scale choice of pdf's ambiguous. Alternatives range from  $p_{\perp}^2$  to  $p_{\perp}^2 + Q_1^2 + Q_2^2$ , with preferred one

$$\mu^2 = p_{\perp}^2 \frac{\hat{s} + Q_1^2 + Q_2^2}{\hat{s}} \quad \left( \sim \frac{-\hat{t}\hat{u}}{\hat{t} + \hat{u}} \right).$$

So far no resolved longitudinal photon, except by multiplicative factor

$$R(\mu^2, Q^2) = 1 + a \frac{4\mu^2 Q^2}{(\mu^2 + Q^2)^2}.$$

## Other model aspects

Initial-state shower cut-off: normally  $Q_0^{\text{sh}} = 1 \text{ GeV}$ .  
 If vector meson/anomalous state 'resolution scale'  $\tilde{Q}_0/k$  is above this, shower cut-off is increased correspondingly.

Final-state radiation: the 'beam remnant' quark in an anomalous branching  $\gamma^* \rightarrow q\bar{q}$  is allowed to radiate, from a scale  $k$  to  $m_0^{\text{sh}} = 1 \text{ GeV}$ , while a VMD beam remnant cannot.

Primordial  $k_{\perp}$ : Gaussian 'small' for VMD,  
 $\sim dk_{\perp}^2/k_{\perp}^2$  for anomalous.

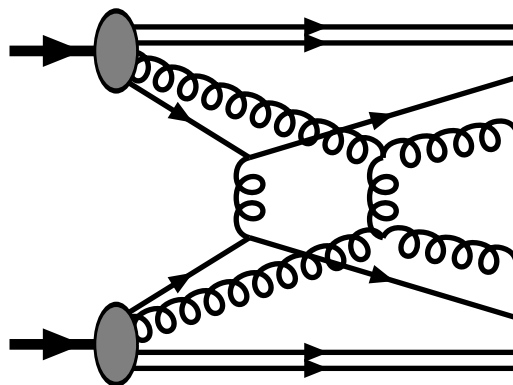
Multiple parton-parton interactions:  
 only affects VMD  $\times$  VMD.

$$\langle n_{\text{int}} \rangle \sim \frac{1}{\sigma_{\text{tot}}} \int_{p_{\perp \text{min}}^2}^{s/4} \frac{d\sigma_{\text{jet}}}{dp_{\perp}^2} dp_{\perp}^2$$

Ansatz for  $\gamma^*$ :

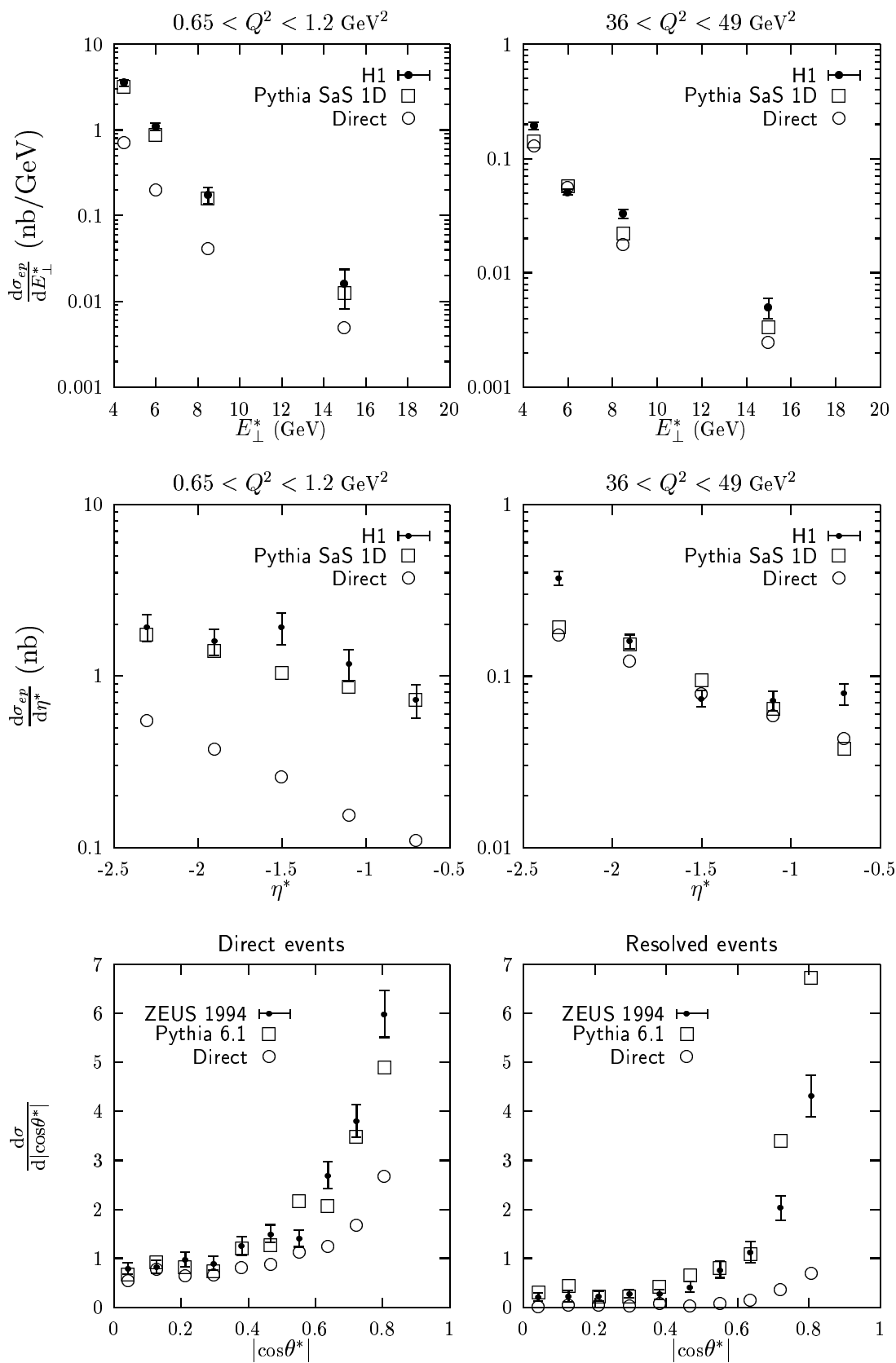
$$\sigma_{\text{tot}} \propto \frac{m_V^2}{m_V^2 + Q^2}$$

$$p_{\perp \text{min}} \propto \sqrt{1 + \frac{Q^2}{m_V^2}}$$



Hadronic data :  $p_{\perp \text{min}} \approx (1.9 \text{ GeV}) \left( \frac{s}{1 \text{ TeV}^2} \right)^{0.08}$

# Comparison with HERA data

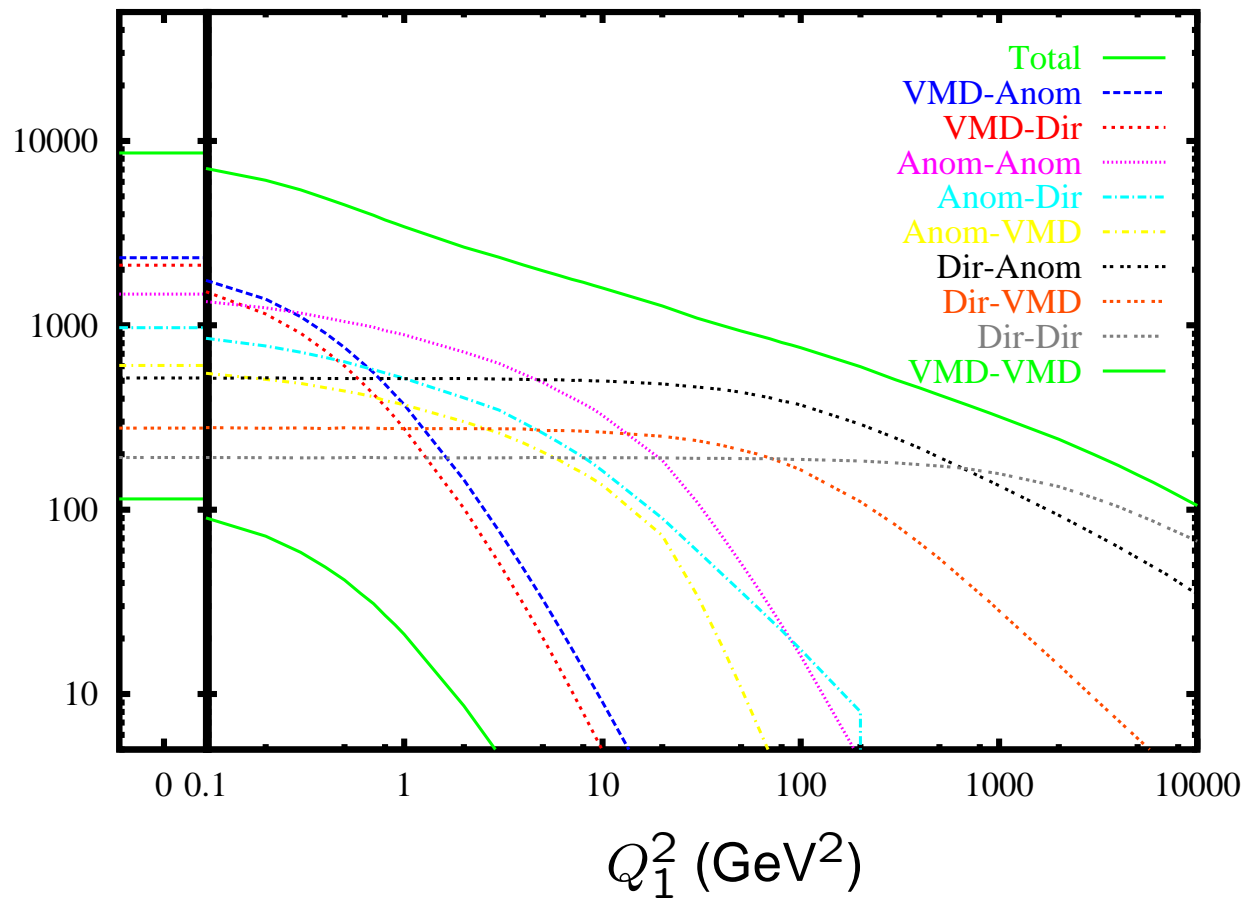




# $e^+e^-$ cross sections

$\sqrt{s_{\gamma^*\gamma^*}} = 100 \text{ GeV}, Q_2^2 = 1 \text{ GeV}^2, p_\perp > 5 \text{ GeV}:$

$\sigma_{\text{jet}}^{\gamma^*\gamma^*}$  (pb)



## Program particulars

To access new  $\gamma^*$  flux convolution:

```
CALL PYINIT('cms', 'gamma/e-', 'gamma/e+', 500D0)
```

Also possible to have  $\gamma^*\gamma^*$  collisions directly:

```
CALL PYINIT('five', 'gamma', 'gamma', 400D0)
```

with P(1,J) and P(2,J) defining momenta and virtualities (P(I,5) < 0 for spacelike ones).

Select jet production with MSEL = 1 (default).

Photon character regulated by MSTP(14):

= **10** : mix direct/VMD/anomalous for real photons;  
D×V = V×D etc. ⇒ 6 classes.

= **20** : (default) mix direct/VMD/anomalous for virtual photons; ⇒ 9 classes.

= **25** : mix direct/resolved for virtual photons; ⇒ 4 classes.

= **other numbers** : individual classes.

**warning** :  $\gamma^*\gamma^* \rightarrow \ell^+\ell^-$  included in dir×dir if not switched off (automatic when mixing).

Possible to specify cuts on  $x_i, y_i, Q_i^2, \theta_i, W^2$  in CKIN(61) - CKIN(78).

Phase space sampled according to

$$\prod_i (dQ_i^2/Q_i^2) (dx_i/x_i) d\varphi_i$$

⇒ full efficiency for  $x_i$  and  $Q_i^2$  cuts.

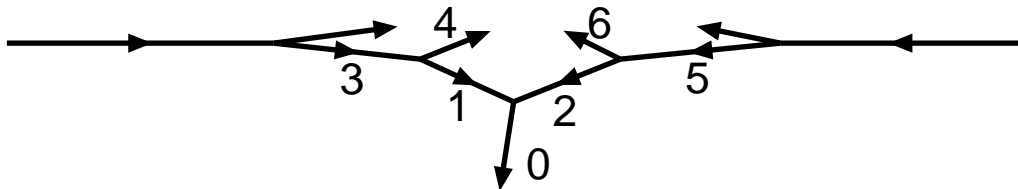
# Photon ISR in $Z^0$ production

(G. Miu & TS, hep-ph/9812455 → PLB)

By-product of study on  $W^\pm$  production in hadron colliders, attempting to combine **matrix-element (ME)** and **parton-shower (PS)** strengths.

Merging strategy: correct hardest emissions in showers so as to reproduce one order higher matrix elements.

$2 \rightarrow 1$  process  $e^+(1) + e^-(2) \rightarrow Z^0(0)$  starting point for backwards shower evolution:



$2 \rightarrow 2$  process  $e^+(3) + e^-(2) \rightarrow \gamma(4) + Z^0(0)$ :

$$\hat{s} = (p_3 + p_2)^2 = \frac{(p_1 + p_2)^2}{z} = \frac{m_Z^2}{z}$$

$$\hat{t} = (p_3 - p_4)^2 = p_1^2 = -Q^2$$

$$\hat{u} = m_Z^2 - \hat{s} - \hat{t} = Q^2 - \frac{1-z}{z} m_Z^2$$

Relate **ME** and **PS** rates:

$$\frac{d\hat{\sigma}}{d\hat{t}} \Big|_{\text{ME}} = \frac{\sigma_0 \alpha_{\text{em}}}{\hat{s} 2\pi} \frac{\hat{t}^2 + \hat{u}^2 + 2m_Z^2 \hat{s}}{\hat{t}\hat{u}}$$

$$\xrightarrow{Q^2 \rightarrow 0} \sigma_0 \frac{\alpha_{\text{em}}}{2\pi} \frac{1+z^2}{1-z} \frac{1}{Q^2} = \frac{d\hat{\sigma}}{dQ^2} \Big|_{\text{PS1}}$$

$$\frac{d\hat{\sigma}}{d\hat{t}} \Big|_{\text{PS1}} = \frac{\sigma_0 \alpha_{\text{em}}}{\hat{s} 2\pi} \frac{\hat{s}^2 + m_Z^4}{\hat{t}(\hat{t} + \hat{u})}$$

Add mirror  $e^+(1) + e^-(5) \rightarrow \gamma(6) + Z^0(0)$ :

$$\frac{d\hat{\sigma}}{d\hat{t}} \Big|_{\text{PS}} = \frac{d\hat{\sigma}}{d\hat{t}} \Big|_{\text{PS1}} + \frac{d\hat{\sigma}}{d\hat{t}} \Big|_{\text{PS2}} = \frac{\sigma_0 \alpha_{\text{em}}}{\hat{s} 2\pi} \frac{\hat{s}^2 + m_Z^4}{\hat{t}\hat{u}}$$

$$R_{ee \rightarrow \gamma Z}(\hat{s}, \hat{t}) = \frac{(d\hat{\sigma}/d\hat{t})_{\text{ME}}}{(d\hat{\sigma}/d\hat{t})_{\text{PS}}} = \frac{\hat{t}^2 + \hat{u}^2 + 2m_Z^2 \hat{s}}{\hat{s}^2 + m_Z^4}$$

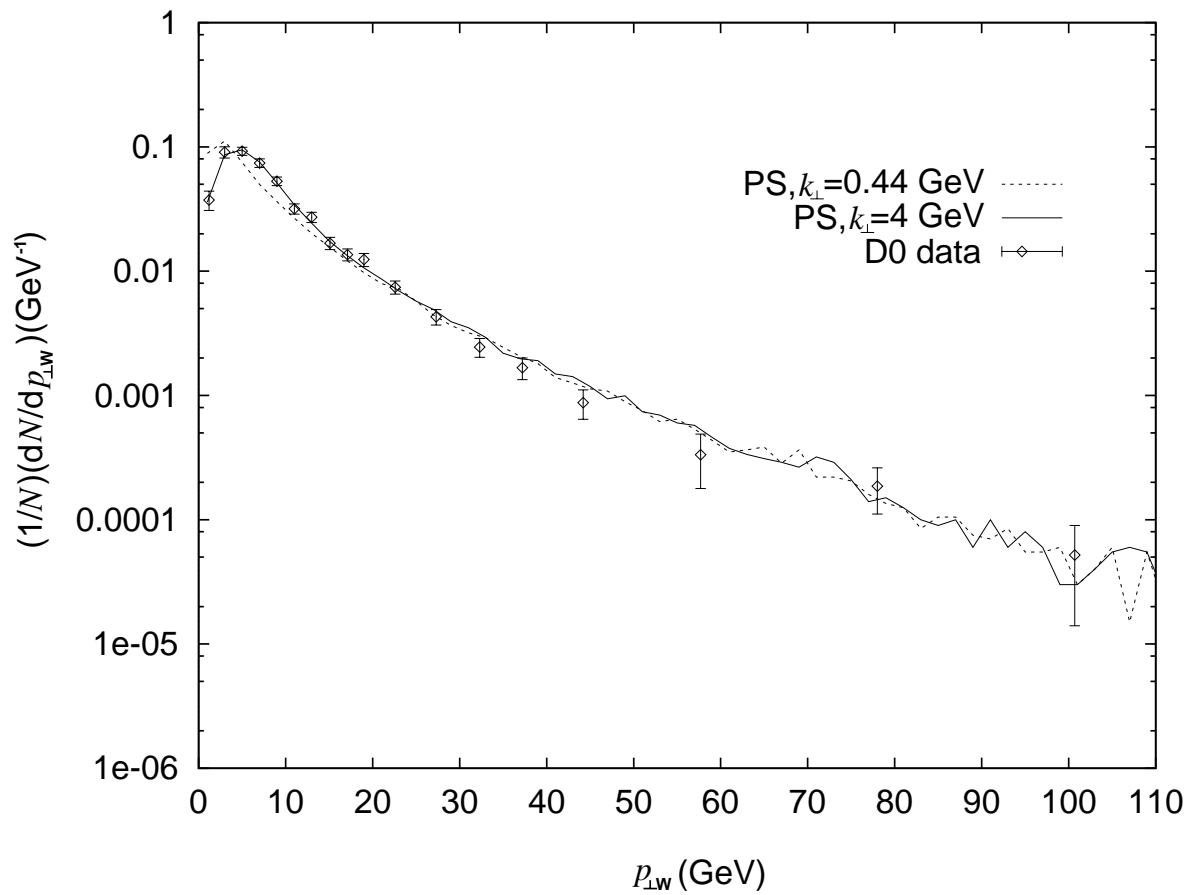
$$\frac{1}{2} < R_{ee \rightarrow \gamma Z}(\hat{s}, \hat{t}) \leq 1$$

Improve **PS**:

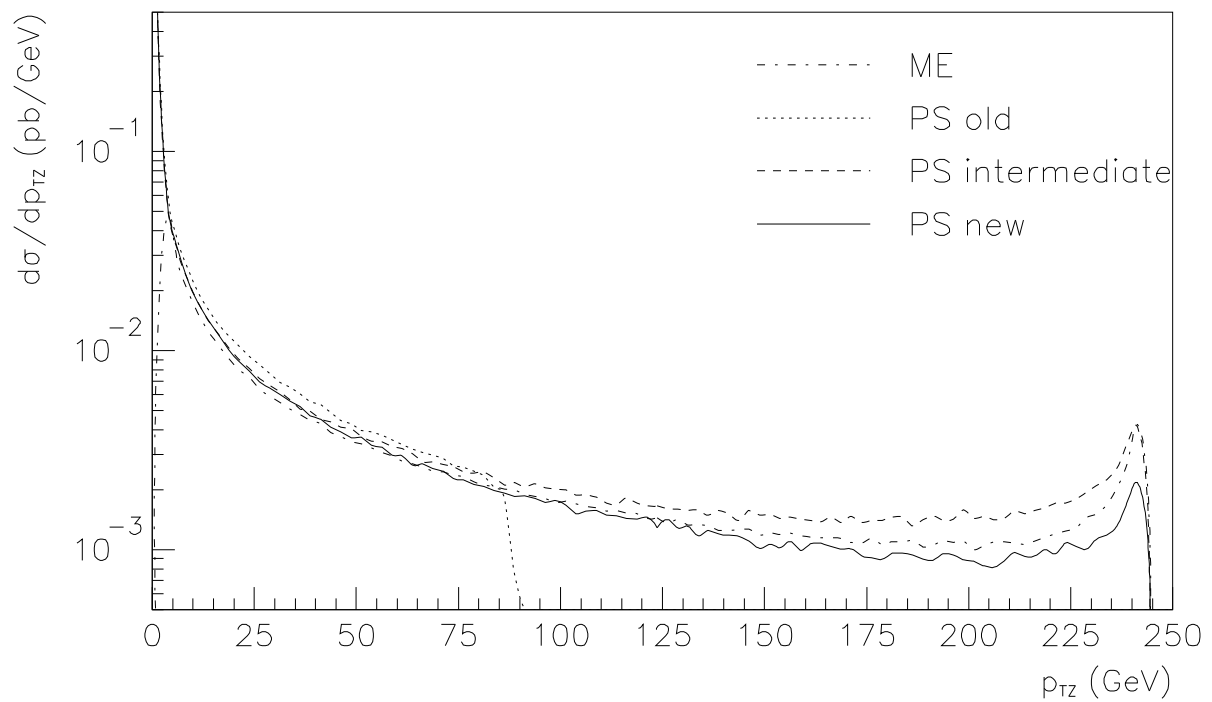
- $Q_{\text{max}}^2 = s$ , not  $Q_{\text{max}}^2 \approx m_Z^2$  (intermediate)
- MC correction by  $R(\hat{s}, \hat{t})$  for first ( $\approx$  hardest) emission on each side (new)

Now default.

$p\bar{p} \rightarrow W^\pm$  at 1.8 TeV:



$e^+e^- \rightarrow Z^0$  at 500 GeV,  
 $80 \text{ GeV} < \sqrt{s} < 100 \text{ GeV}$ :



# Doubly-charged Higgses

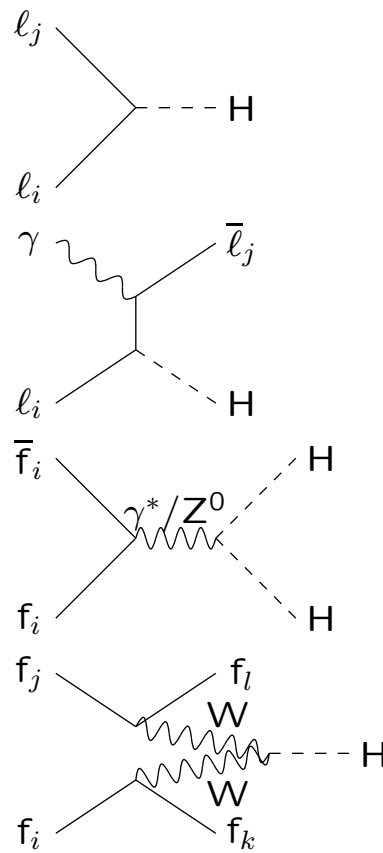
(K. Huitu, J. Maalampi, A. Pietilä, M Raidal, private communication; G. Barenboim, KH, JM, MR, PLB394 (1997) 132)

Based on left–right symmetric scenario.

New particles:  $\nu_{iR}$ ,  $W_R^\pm$ ,  $H_L^{\pm\pm}$ ,  $H_R^{\pm\pm}$ .

New processes:

- 341  $l_i l_j \rightarrow H_L^{\pm\pm}$
- 342  $l_i l_j \rightarrow H_R^{\pm\pm}$
- 343  $l_i^\pm \gamma \rightarrow H_L^{\pm\pm} e^\mp$
- 344  $l_i^\pm \gamma \rightarrow H_R^{\pm\pm} e^\mp$
- 345  $l_i^\pm \gamma \rightarrow H_L^{\pm\pm} \mu^\mp$
- 346  $l_i^\pm \gamma \rightarrow H_R^{\pm\pm} \mu^\mp$
- 347  $l_i^\pm \gamma \rightarrow H_L^{\pm\pm} \tau^\mp$
- 348  $l_i^\pm \gamma \rightarrow H_R^{\pm\pm} \tau^\mp$
- 349  $f_i \bar{f}_i \rightarrow H_L^{++} H_L^{--}$
- 350  $f_i \bar{f}_i \rightarrow H_R^{++} H_R^{--}$
- 351  $f_i f_j \rightarrow f_k f_l H_L^{\pm\pm}$
- 352  $f_i f_j \rightarrow f_k f_l H_R^{\pm\pm}$



Typical decays:

- $H_L^{++} \rightarrow l_i^+ l_j^+, W_L^+ W_L^+$
- $H_R^{++} \rightarrow l_i^+ l_j^+, W_R^+ W_R^+$
- $W_R^+ \rightarrow q \bar{q}', l_i^+ \nu_{lR}$

Status: working, but still a few factors of  $\sim 2$  to sort out.

# Other news

2–, 4– and 6–fermion standard interfaces for showers and hadronization:

```
CALL PY2FRM(IRAD, ITAU, ICOM)
```

```
CALL PY4FRM(ATOTSQ, A1SQ, A2SQ, ISTRAT,  
&IRAD, ITAU, ICOM)
```

```
CALL PY6FRM(P12, P13, P21, P23, P31, P32, PTOP,  
&IRAD, ITAU, ICOM)
```

$P_{ij}$  : relative probability that first (second) fermion is paired with  $i$ 'th ( $j$ 'th) antifermion.

PTOP : absolute probability for  $t\bar{t}$  event. If  $t\bar{t}$  is selected, the  $P_{ij}$  are not used. The  $b\bar{b}$  must be first fermion pair.

Process 36,  $e\gamma \rightarrow \nu W (\Rightarrow ee \rightarrow e\nu W)$  :  
W decay angle ME now included.

New function PYMRUN for running  $\overline{MS}$  masses, e.g. in Higgs production  $\Rightarrow$  PMAS(KC,1) free to use for “on-shell” masses (e.g. for charm production rates).

Technicolour processes upgraded;  
more to come.

New processes specifically for  $\tilde{b}$  production.

Many bugs fixed, and (no doubt) new ones introduced.

# On to C++!

(L. Lönnblad, hep-ph/9810208 → CPC; M. Bertini, TS)

## Why Fortran → C++?

- SLAC →, FNAL →, CERN → LHC era.
- Industrial standard.
- Educational and professional continuity for students.
- Better to program — for experts.
- User-friendly interfaces — for the rest of us (cooperation with GEANT4, LHC++).

## Milestones:

- January 1998: project formally started (Leif ~half-time).
- Exists today: strategy document, code for the event record and the particle object.
- In progress: particle data and other data base handling, event generation handler structure, string fragmentation.
- Next: decay routines, very simple matrix elements.
- By summer: proof of concept.
- End 2000: most of current PYTHIA functionality (?).
- ??: more and better than current PYTHIA.

Input welcome: `leif@thep.lu.se`